

# PROLOGUE

**The Geometrization of Thermodynamics**  
(particles via Caratheodory)

VS

**The Topolization of Thermodynamics**  
(particles and waves via Grassmann)

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**Warning**

**The Topology of Particles**

is not the same as

**The Topology of Waves**

September 22, 2012

# PROLOGUE

**Thermodynamic states  $\Leftrightarrow$  topological structures.**

**Equilibrium states  $\Leftrightarrow$  Pfaff Topological Dimension  $M \leq 2$**

**Non-Equilibrium states  $\Leftrightarrow$  Pfaff Topological Dimension  $M \geq 3$**

**Homotopic methods acting on thermodynamic topological structures can mathematically distinguish continuous topological evolution of **Irreversible, Non-Deterministic, Non-Integrable, or Self-Organizing Emergent processes****

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# PROLOGUE

**An exterior differential 1-form,  $\mathbf{A}$ , defines a thermodynamic state of  $k$  geometric functions and differentials.**

**The neighborhood topology defined by  $\mathbf{A}$  is composed of  $\mathbf{M}$  functions and differentials **such that  $\mathbf{M} \leq k$ .****

**$\mathbf{M}$  is defined as the Pfaff Topological Dimension of  $\mathbf{A}$   
( $\mathbf{M}$  is easy to compute in terms of the Pfaff Sequence)**

## The Pfaff Sequence and the Pfaff Topological Neighborhood Dimension

The Geometrical 1-form  $\mathbf{A} = x dx - y dy + (k dz - \omega dt)$  has 4 functions, 2 constants and 4 differentials

But

$$\mathbf{PS}(\mathbf{A}) = [\mathbf{A}, d\mathbf{A} = 0, \mathbf{A} \wedge d\mathbf{A} = 0, 0] = [\mathbf{A}, 0, 0, 0].$$

(The Pfaff Sequence has only 1 non-zero entry).

Hence the Pfaff Topological Dimension =  $\mathbf{PTD}(\mathbf{A}) = \mathbf{M} = 1$ .

The topological neighborhood environment is an **equilibrium** thermodynamic state.

The topological minimum Pfaffian =  $\mathbf{A} = d\Psi$  has 1 differential  
 $\Psi = \{(x^2 - y^2)/2 + (kz - \omega t)\}$

September 22, 2012

# PROLOGUE

This essay is a summary of work that was started in 1962 and led to topological non-equilibrium properties that include:

Topological Spin in non-equilibrium systems (1971),  
Topological Torsion in turbulent systems (1977),  
Topological Emergence of Falaco Solitons (1987)  
Homotopic Continuous Topological Evolution (1987-1991),  
Thermodynamic Irreversibility and the Arrow of Time (2003),  
Topological Torsion and Spin in Plasmas (2004),  
Thermodynamic Topologies from Pfaffian 1-forms (2009),  
Thermodynamic Topologies from Current N-1-forms (2010)  
Topologies of Waves versus Topologies of Particles (2012)

# PROLOGUE

## Particles vs. Waves

The **T0** Topology of **Particles**  
is not the same as  
the **Not T0** Topology for **waves**,

**but the two topologies can coexist**

**Physical measurements** are over **finite** space-time intervals.

**Finite topologies** are of 2 types: **T0** and **Not-T0**.

The **T0** topology should be used for systems of **Particles**.

The **Not-T0** topology should be used for systems of **Waves**.

The **T0** topologies can be generated from the functional properties of an exterior differential 1-form, **A**, of **Action-potentials**, creating a **Particle** topology with **uniquely distinguishable** singlet subsets.

The **Not-T0** topologies can be generated in terms of the properties of an  $N-1$  form, **[C,  $\rho$ ]**, or current density, representing a **Complex-Wavelet, or ensemble**, topology, where **all** of the subsets are **NOT uniquely distinguishable**.



The complete lattice structure for an extremal **Non T0** topology of 3 ingredients without closure axioms. All subsets have the same closure and cannot be distinguished. There are 4 different Non T0 topologies of 3 ingredients.

$$X := \{a, b, c\}$$

$$LS := \{\{\}, \{a, b, c\}\}$$

Not-T0 N=3 Example 1, Indiscrete, Connected, as all sub sets are dense,  $\text{Bnd of Bnd}(S) = \{\}$   
 Is LS a topology = true, connected = true, Kolmogorov.T0 = false, Hausdorff.T2 = false

COMPLETE Lattice Structure

Subset S	Int(S)	Ext(S)	Bnd(S)	Clo(S)	Lim(S)	IsoClo(S)	IsoCar(S)
{a}	{}	{}	{a, b, c}	{a, b, c}	{b, c}	{a}	{}
{b}	{}	{}	{a, b, c}	{a, b, c}	{a, c}	{b}	{}
{c}	{}	{}	{a, b, c}	{a, b, c}	{a, b}	{c}	{}
{a, b}	{}	{}	{a, b, c}	{a, b, c}	{a, b, c}	{}	{a, b}
{a, c}	{}	{}	{a, b, c}	{a, b, c}	{a, b, c}	{}	{a, c}
{b, c}	{}	{}	{a, b, c}	{a, b, c}	{a, b, c}	{}	{b, c}
{a, b, c}	{a, b, c}	{}	{}	{a, b, c}	{a, b, c}	{}	{a, b, c}

Closed-Open subsets of LS are =  $\{\{\}, \{a, b, c\}\}$

Some singletons have the same closure and cannot be distinguished  
 The {1,2,3} array of CLOSURE elements is =  $(\{a, b, c\}, \{a, b, c\}, \{a, b, c\})$

The complete lattice structure for an extremal **T0** topology of 3 ingredients with the Kolmogorov closure axioms. The singlet subsets have distinguishable closures. There are 5 different Kolmogorov T0 topologies of 3 ingredients.

$X := \{a, b, c\}$   
 $LS := \{\{\}, \{a\}, \{b\}, \{c\}, \{a, b\}, \{a, c\}, \{b, c\}, \{a, b, c\}\}$   
 $N=3$  T2 poset 1 Discrete, Bnd of Bnd(S) =  $\{\}$  as all Bnd(S) =  $\{\}$

Is LS a topology = true, connected = false, Kolmogorov.T0 = true, Hausdorff.T2 = true

COMPLETE Lattice Structure

Subset S	Int(S)	Ext(S)	Bnd(S)	Clo(S)	Lim(S)	IsoClo(S)	IsoCar(S)
{a}	{a}	{b, c}	{}	{a}	{}	{a}.Seg	{}
{b}	{b}	{a, c}	{}	{b}	{}	{b}.Seg	{}
{c}	{c}	{a, b}	{}	{c}	{}	{c}.Seg	{}
{a, b}	{a, b}	{c}	{}	{a, b}	{}	{a, b}.Seg	{}
{a, c}	{a, c}	{b}	{}	{a, c}	{}	{a, c}.Seg	{}
{b, c}	{b, c}	{a}	{}	{b, c}	{}	{b, c}.Seg	{}
{a, b, c}	{a, b, c}	{}	{}	{a, b, c}	{}	{a, b, c}.Seg	{}

Closed-Open subsets of LS are =  $\{\{\}, \{a\}, \{b\}, \{c\}, \{a, b\}, \{a, c\}, \{b, c\}, \{a, b, c\}\}$   
The {1,2,3} array of CLOSURE elements is =  $\{\{a\}, \{b\}, \{c\}\}$

This essay will describe both  
equilibrium and non-equilibrium thermodynamics  
from the basis of Continuous Topological Evolution,  
resulting in

## **A UNIVERSAL Category Theory of Topological Thermodynamics**

It is remarkable that the method introduces a  
topological formalism useful to the description of  
irreversible, non-deterministic dynamics  
of both **Particles** and **Waves**,  
in terms of **only two types** of finite coexistent topologies.

A Marriage between **Particles** and **Waves**  
requires the use of the two different topologies.

For example, a **gravitational** metric theory  
(of **particles**),  
cannot be merged with a **quantum theory**  
(of **waves**),  
without the simultaneous use of the  
two finite different topologies,  
**T0** and **Not T0**.

Geometric metric curvature is not a  
Topological invariant

# Basic Ideas 1.

1. The topology of **Particles** is not the same as the topology of **Waves**.
2. Only the discrete topology (Hausdorff T2) of **Particles** is metrizable. *Relativity theorists beware.*
3. The Kuratowski **T0 poset 3 Particle** topology is **NOT metrizable**.

Hence a theory of Gravity based on metric and distinguishable particles is not topologically complete.

Is the gravity of **HOT** bodies  
different than the gravity of **COLD** bodies?

## Basic Ideas 2.

4. The indiscrete **Not T0 Wave** topologies are not metrizable.

5. The **particle** topologies can coexist with the **wave** topologies.

6. Topological evolution can be described by the action of the Lie differential **HOMOTOPIC** Operator acting on differential forms.

The **Homotopy** of a particle topology based on a 1-form of Action is equivalent to the universal **Topological First Law of Thermodynamics**.

The **Homotopic** differential becomes the Covariant differential when the process is adiabatic.

# Applications of The Category Theory of Topological Thermodynamics

This work was motivated more than 40 years ago by the challenge of the Clay Institute regarding the properties of the Navier-Stokes equations and their relationship to hydrodynamic turbulence.

To replicate a statement made by the Clay Institute:

*"The challenge is to make substantial progress toward a mathematical theory which will unlock the secrets hidden in the Navier-Stokes equations."*

**Applications of**  
**The Category Theory of**  
**Topological Thermodynamics**

**can demonstrate**

**Irreversible Continuous**  
**Topological Evolution**



**Applications of**  
**The Category Theory of**  
**Topological Thermodynamics**

**can prove that the**

**Navier-Stokes equations have**  
**Turbulent solutions**

<http://www22.pair.com/csdc/pdf/topturb.pdf>

The **Category theory of Topological Thermodynamics** and any exterior differential 1-form,  $\mathbf{A}$ , of rank 4, can be used to

1. generate a disconnected Cartan Topology that defines a **non-equilibrium** thermodynamic system of particles.

Use the Grassmann algebra of Pfaffians to define

$\mathbf{A} =$  Topological 1-form of Action

$\mathbf{F} = d\mathbf{A} =$  Topological 2-form of Vorticity

$\mathbf{H} = \mathbf{A} \wedge \mathbf{F} =$  Topological 3-form of Torsion

$\mathbf{K} = \mathbf{F} \wedge \mathbf{F} =$  Topological 4-form of Parity

Then use  $\mathbf{A}, \mathbf{F}, \mathbf{H}, \mathbf{K}$  to construct a topology.

Kuratowski Topology  $T_4 \cong$  Kolmogorov  $T_0$ , poset 3

From any 1-form in 4D:  $A = A_k(x^m)dx^k$

Pfaff Sequence =  $\{A, F = dA, H = A \wedge F, K = F \wedge F\}$

Basis subsets  $\{A, K_{Cl}(A), H, K_{Cl}(H)\} = \{A, A \cup F, H, H \cup K\}$

T4{Lattice Structure  $\{X, \emptyset, A, H, A \cup F, H \cup K, A \cup H, A \cup H \cup K, A \cup F \cup H\}$

Complete Lattice Structure

Subset	LimPt(S)	Int(S)	Ext(S)	Bnd(S)	Clos(S)
$\emptyset$	$\emptyset$	$\emptyset$	$X$	$\emptyset$	$\emptyset$
$A$	$F$	$A$	$H \cup K$	$F$	$A \cup F$
$F$	$\emptyset$	$\emptyset$	$A \cup H \cup K$	$F$	$F$
$H$	$K$	$H$	$A \cup F$	$K$	$H \cup K$
$K$	$\emptyset$	$\emptyset$	$A \cup F \cup H$	$K$	$K$
$A \cup F$	$F$	$A \cup F$	$H \cup K$	$\emptyset$	$A \cup F$
$A \cup H$	$F, K$	$A \cup H$	$\emptyset$	$F \cup K$	$X$
$A \cup K$	$F$	$A$	$H$	$F \cup K$	$A \cup F \cup K$
$F \cup H$	$K$	$H$	$A$	$F \cup K$	$F \cup H \cup K$
$F \cup K$	$\emptyset$	$\emptyset$	$A \cup H$	$F \cup K$	$F \cup K$
$H \cup K$	$K$	$H \cup K$	$A \cup F$	$\emptyset$	$H \cup K$
$A \cup F \cup H$	$F, K$	$A \cup F \cup H$	$\emptyset$	$K$	$X$
$F \cup H \cup K$	$K$	$H \cup K$	$A$	$F$	$F \cup H \cup K$
$A \cup H \cup K$	$F, K$	$A \cup H \cup K$	$\emptyset$	$F$	$X$
$A \cup F \cup K$	$F$	$A \cup F$	$H$	$K$	$A \cup F \cup K$
$X$	$F, K$	$X$	$\emptyset$	$\emptyset$	$X$

The **Category theory of Topological Thermodynamics** and any exterior differential 1-form,  $\mathbf{A}$ , of rank 4, can be used to

1. generate a disconnected Cartan Topology that defines a **non-equilibrium** thermodynamic system of particles.
2. generate a Jacobian correlation matrix,  $[\partial \mathbf{A}_k / \partial x^k]$  that has a singular characteristic set of rank 3, which is a morphism of a **universal van der Waals gas**.
3. generate a unique process current of **Topological Torsion**,  $\mathbf{T} = \mathbf{A} \wedge d\mathbf{A}$ , which describes an irreversible thermodynamic process in a non-equilibrium system.

Topological Torsion  $\int \mathbf{A} \wedge d\mathbf{A}$  is a key design tool for  
controlling and understanding

**Dissipative Structures and**

**TURBULENCE**

Almost **NO** Engineers and Very Few Physicists Understand

**TOPOLOGICAL TORSION**

(pity)



*As his students see him*

**By Professor R. M. Kiehn**

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## Prologue

I am not an **algebraist**.

I am not a **topologist**.

I am an **applied engineering physicist**, looking for a better, **non-phenomenological**, universal way to understand **non-equilibrium** thermodynamic systems and irreversible processes.

The goal is to find useful (and practical) applications of the universal topological methods in many different disciplines of science and engineering.

Starting in 1962, after leaving Los Alamos, the objective of my research has been to formulate an extension of tensor and variational techniques that would yield a **universal** description of **non-equilibrium** systems and **irreversible turbulent** processes, applicable to both particle and statistical physical disciplines.

The result has been the creation of a **category theory of topological thermodynamics**, which includes both **particle** and **statistical** features as **distinct topological structures**. The method has produced useful patents, a better understanding of turbulence in fluids and plasmas, and a proof of the Prigogine conjecture of emergence of metastable states far from equilibrium (in both quantum and macroscopic disciplines).



Mathematicians define a **Category Theory, C**, to be

1. A collection of abstract **objects**,  $\text{ob}(\mathbf{C})$ , with
2. **morphisms** that map an object **X** to another object **Y**.

The **Category of Topological Thermodynamics** is

1. A collection of the **2** finite types of **topological spaces** which are defined by the exterior differential 1-form, **A**, or the N-1 form current density, **C, ρ** and which are used to represent the **2 thermodynamic systems** subject to
2. (possibly **non-invertible but continuous**) **Homotopic morphisms** that map exterior differential forms into other exterior differential forms, representing **processes**, which include **topological or geometrical change**.

**Theorem 1:** Any exterior differential 1-form of Action  $\mathbf{A}$  can be used to generate the topological structure suitable to describe thermodynamic systems of distinguishable particles.

Using the Grassmann algebra compute the Pfaff Sequence of  $\mathbf{A}$ ,

$$\text{PS}(\mathbf{A}) = [\mathbf{A}, d\mathbf{A}, \mathbf{A} \wedge d\mathbf{A}, d\mathbf{A} \wedge d\mathbf{A} \dots]$$

$\mathbf{A}$  = Topological Action 1-form

$d\mathbf{A}$  = Topological Vorticity 2-form

$\mathbf{A} \wedge d\mathbf{A}$  = Topological Torsion 3-form

$d\mathbf{A} \wedge d\mathbf{A}$  = Topological Parity 4-form

**Theorem 2:** The minimum number  $M$  of functions required to define the topological neighborhood of  $A$  defines the Pfaff Topological Dimension,  $PTD(A)$ , of the 1-form  $A$ .  $M$  is less than or equal to the geometric dimension,  $k$ , of  $A$ .

The number of non-zero entries in the Pfaff Sequence  
Determines the Pfaff Topological Dimension of  $A$ .

$$PS = [A, 0, 0, 0] \quad \Rightarrow \quad PTD(A) = 1$$

$$PS = [A, dA, 0, 0] \quad \Rightarrow \quad PTD(A) = 2$$

$$PS = [A, dA, A \wedge dA, 0] \quad \Rightarrow \quad PTD(A) = 3$$

$$PS = [A, dA, A \wedge dA, dA \wedge dA] \quad \Rightarrow \quad PTD(A) = 4$$

Define the top Environment as a  $\text{PTD}(\mathbf{A}) = 4$  system

**Equilibrium**  $\text{PTD}(\mathbf{A}) = 1$  subsystems do **not** exchange heat or mass with the environment

**Isolated Equilibrium**  $\text{PTD}(\mathbf{A}) = 2$  subsystems do **not** exchange heat or mass with the environment

**NON-Equilibrium Closed**  $\text{PTD}(\mathbf{A}) = 3$  subsystems can exchange **heat**, but not **mass**, (**or mass**, but not **heat**), with the environment

**NON-Equilibrium Open**  $\text{PTD}(\mathbf{A}) = 4$  subsystems can exchange both **heat** and **mass** with the environment

Thermodynamic **processes J** are used to define topologically continuous, differentiable, morphisms (**maps**) between the initial state **A** and the final state, **Q**.

$$\text{Map } \mathbf{A} \Rightarrow \mathbf{Q}$$

The process direction field, **J** can be composed of **vectors** and macroscopic **spinors**

A Homotopy =  $\mathbf{H} \mathbf{d} + \mathbf{d} \mathbf{H}$  can be formulated in terms of the Lie differential  $L_{(\mathbf{J})}$  with respect to a current, **J**:

$$L_{(\mathbf{J})} = \mathbf{i}(\mathbf{J}) \mathbf{d} + \mathbf{d} (\mathbf{i}(\mathbf{J})).$$

If applied to differential forms, the definitions lead to:

## Cartan's Magic Homotopy Formula 1899-1923:

$$L_{(\mathbf{J})} \mathbf{A} = i(\mathbf{J})d\mathbf{A} + d(i(\mathbf{J})\mathbf{A})$$

**Theorem:** The homotopy operator,  $L_{(\mathbf{J})}\mathbf{A}$ , Cartan, operating on exterior differential forms that are C2 smooth (C2 differentiable)

**is topologically continuous.**

The 4 components of  $\mathbf{C}$  can consist of both vector/spinor **direction fields** ( $\sim$  flow) and can be used to define a 3-form Current density.

The Cartan homotopy maps the Limit points,  $d\Sigma$ , of the ingredients,  $\Sigma$ , of the **initial** topology into the Closure,  $\Sigma \cup d\Sigma$ , of the ingredients of the **final** topology

For 1-forms, the initial state topology of **A** is not necessarily the same as the final state topology of **Q** !!!

$$L_{(J)} A = i(J)dA + d(i(J)A) \Rightarrow Q$$

**In fact, Topological Evolution is a necessary requirement for process irreversibility.**

*The fact that any differential 1-form can be used to define a topological space came to my attention about 1987, after some 20 years of studying Cartan's utilization of differential forms and their application to the theory of integral invariants. Early on I defined this topological space as the disconnected Cartan topology of exterior differential forms.*

*However, it was not until 2009, when I attended (in attempt to learn more about modern topology) my first conference for professional topologists at Hacetepe, Turkey, that I realized that the Cartan topology of 4 differential form ingredients was exactly equivalent to a Kuratowski representation of the disconnected Kolmogorov  $T_0$  poset 3 topology, based upon the power set of 4 exterior differential forms as ingredients.*



# Sequence of Ideas 1987-2009

Following the lead of E. Cartan, the "laws of motion" that describe the continuous evolution of the T0 "particle" topologies will be generated by applying the homotopic Lie differential (not derivative) to the exterior differential 1-form of Action,  $\mathbf{A}$ , that encodes the thermodynamic system..

# Sequence of Ideas 1987-2009

The result is to homotopically map  $\mathbf{A} \Rightarrow \mathbf{Q}$

$$\mathbf{L}_{(\mathbf{J})} \mathbf{A} = \mathbf{i}(\mathbf{J})d\mathbf{A} + d(\mathbf{i}(\mathbf{J})\mathbf{A}) \Rightarrow \mathbf{Q}$$

By change of notation:

$$\mathbf{W} = \mathbf{i}(\mathbf{J})d\mathbf{A} \quad \text{Work 1-form,}$$

$$\mathbf{U} = \mathbf{i}(\mathbf{J})\mathbf{A} \quad \text{Internal energy,}$$

$$\mathbf{Q} = \mathbf{L}_{(\mathbf{J})} \mathbf{A} \quad \text{Heat 1-form}$$

Cartan's Magic Formula becomes

(for all systems,  $\mathbf{A}$ , and all processes,  $\mathbf{J}$ )

**universally** and **topologically** equivalent to

$$\mathbf{L}_{(\mathbf{J})} \mathbf{A} = \mathbf{W} + \mathbf{d}(\mathbf{U}) = \mathbf{Q}$$

Which is a topological formulation of

**The FIRST LAW of THERMODYNAMICS!!**

Even if the process  $\mathbf{J}$  induces irreversible topological change!

The 1<sup>st</sup> Cartan Magic formula represents a system of PDE's :

$$Q - W = d(U)$$

Hence , the

### FIRST LAW of THERMODYNAMICS

goes beyond DeRham cohomology theory , which requires that both  $Q$  and  $W$  be closed and integrable,  $dQ = dW = 0$ .

The more general “cohomology” statement implies  $Q$  and  $W$  are **not necessarily integrable**, but their limit sets are related and not necessarily zero.

$$Q \wedge dQ = d(U) \wedge dW$$

## RESULTS

DeRham cohomology can be applied to  
**Equilibrium Thermodynamics,**

$$\text{PTD}(\mathbf{A}) < 3$$

but the Homotopic Cartan *extended*  
cohomology can be applied to

**NON-Equilibrium Thermodynamics**

$$\text{PTD}(\mathbf{A}) > 2$$

The current, or flow,  $\mathbf{C} = \rho \mathbf{V}_4$ , is defined as a  
“differential direction field”  
that is NOT necessarily a derivative,  
nor a generator of a 1-parameter group.

$$\mathbf{C} = \rho \mathbf{V}_4 = [\mathbf{C}^x, \mathbf{C}^y, \mathbf{C}^z, \mathbf{C}^t]$$

On a 4D variety, the contraction of  $\mathbf{C}$  with a 4-form differential volume element yields a 3-form Current density of 4 coefficients which may or may not have a Zero divergence. The direction field,  $\mathbf{C} = \rho \mathbf{V}_4$ , can be an element of a differential semi-group, or an element of a complex macroscopic spinor space.

# Additional Ideas 1987-2009

## Macroscopic Spinors $S_k$

Any **anti-symmetric matrix**, such as the 2-form  $F = dA$ ,  
will have complex direction field eigenvectors,  $S_k$ ,  
with zero norm (length) but finite area

$\langle S_k S_k \rangle$  equals 0, but  $\langle S_k S_j \rangle$  does not equal 0.

# Additional Ideas 1987-2009

## Macroscopic Spinors of $F=dA$

The Work 1-form  $W$  is defined in terms process ,  $\rho V_4$ ,  
and the 2-form  $F = dA$ .

$$W = i(\rho V_4)F$$

The eigen values of the antisymmetric matrix  $F$  are  
Two imaginary conjugate pairs, if  $PTD(A) = 4$ ,  
1 imaginary conjugate pair, and one 0, if  $PTD(A) = 3$ ,  
1 imaginary conjugate pair, if  $PTD(A) = 2$ ,



# Additional Ideas 1987-2009

**Macroscopic Spinors are**

**Majorana Spinors !**

# Additional Ideas 1987-2009

The correlation matrix,  $[\partial \mathbf{A}_k / \partial \mathbf{x}^k]$ , defines a Caley-Hamilton characteristic polynomial,  $\Theta$ , of 4th degree in terms of similarity invariants. The hypersurface

$$\Theta = 0,$$

can be interpreted as a **thermodynamic phase function** in the symplectic topological space  $\text{PTD}(\mathbf{A}) = 4$ ,

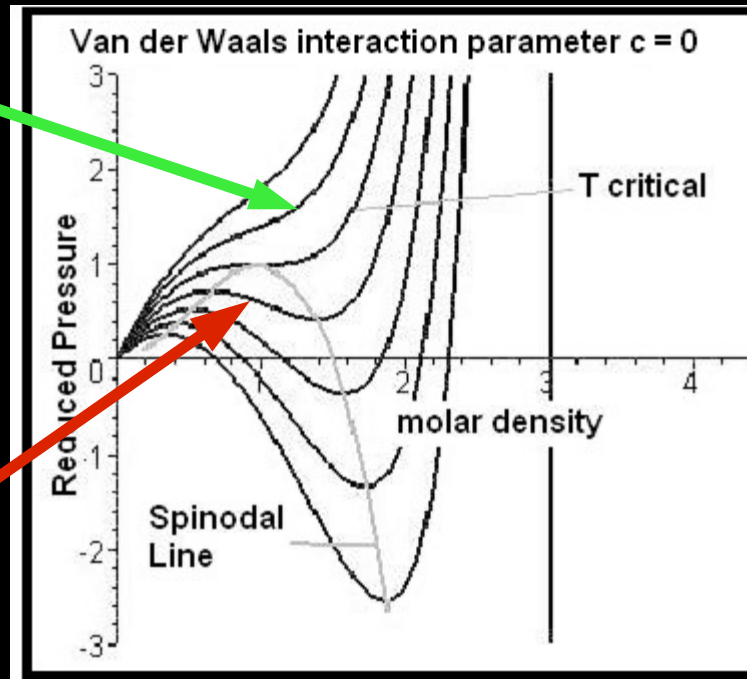
and has the ubiquitous (homoeomorphic) properties of a **van der Walls gas** in the contact topological space,  $\text{PTD}(\mathbf{A}) = 3$ .

These particle aspects have been well documented with many examples in numerous publications over the years, starting in 1962. See <http://www.cartan.pair.com>

# Additional Ideas 1987-2009

Complex eigenvectors represent anti-symmetric features of the Jacobian matrix for a **van der Waals gas**.

Above the critical isotherm there is 1 real eigenvector and 2 complex Macroscopic Spinor eigenvectors of zero norm.



Below  $T$  critical, there are 3 real eigenvectors.

# Additional Ideas 1987-2009

Any 1-form of Action,  $A$ , can be used to define a unique non-zero 3-form process,

$$i(T_4)\Omega = i(T_4) dx^{\wedge}dy^{\wedge}dz^{\wedge}dt = A^{\wedge}dA,$$

with the properties such that

$$i(T_4)A = 0 = i(T_4)Q \text{ (transversely adiabatic),}$$

$$i(T_4)dA = W = \sigma A = Q,$$

$T_4$  REPRESENTS THE 4 COMPONENTS OF

**TOPOLOGICAL TORSION**

# Additional Ideas 1987-2009

A projectively unique current 3-form is deduced from the functional properties that define the thermodynamic system.

$$i(T_4)\Omega = i(T_4) dx^{\wedge}dy^{\wedge}dz^{\wedge}dt = \mathbf{A}^{\wedge}d\mathbf{A},$$

$T_4$  REPRESENTS THE 4 COMPONENTS OF  
TOPOLOGICAL TORSION

$T_4$  is an ARTIFACT of  $PTD(\mathbf{A}) = 4$   
and non-equilibrium systems  
(and vanishes for equilibrium systems.

# Topological Torsion Properties

$\mathbf{T}_4$  on  $\Omega_4$  : Properties of Topological Torsion

$$i(\mathbf{T}_4)\Omega_4 = i(\mathbf{T}_4)dx^{\wedge}dy^{\wedge}dz^{\wedge}dt = A^{\wedge}dA,$$

$$i(\mathbf{T}_4)i(\mathbf{T}_4)\Omega_4 = 0, \text{ which implies that}$$

$$\text{Work 1-form } W = i(\mathbf{T}_4)dA = \sigma A,$$

$$dW = d\sigma^{\wedge}A + \sigma dA = dQ,$$

$$\text{Internal Energy } U = i(\mathbf{T}_4)A = 0, \quad \mathbf{T}_4 \text{ is associative,}$$

$$i(\mathbf{T}_4)dU = 0$$

$$i(\mathbf{T}_4)Q = 0 \quad \mathbf{T}_4 \text{ is adiabatic}$$

$$L_{(\mathbf{T}_4)}A = \sigma A, \quad \mathbf{T}_4 \text{ is homogeneous and self-similar}$$

$$L_{(\mathbf{T}_4)}dA = d\sigma^{\wedge}A + \sigma dA = dQ,$$

$$Q^{\wedge}dQ = L_{(\mathbf{T}_4)}A^{\wedge}L_{(\mathbf{T}_4)}dA = \sigma^2 A^{\wedge}dA \neq 0, \quad \mathbf{T}_4 \text{ is irreversible,}$$

$$dA^{\wedge}dA = d(A^{\wedge}dA) = d\{(i(\mathbf{T}_4)\Omega_4)\} = (div_4\mathbf{T}_4)\Omega_4,$$

$$L_{(\mathbf{T}_4)}\Omega_4 = d\{(i(\mathbf{T}_4)\Omega_4)\} = (2\sigma)\Omega_4, \quad \mathbf{T}_4 \text{ causes } \Omega_4 \text{ expansion}$$

**The 4 divergence of  $T_4$**

$$d(i(T_4)\Omega) = 4\text{Div}(T_4)\Omega = \sigma \Omega$$

**Determines the Space-Time**

**Irreversible Dissipation**

**Coefficient  $\sigma$**

# And Now some heresy

**Energy is conserved in classical processes of  
Equilibrium thermodynamics.  $\text{PTD}(\mathbf{A}) < 3$**

**But Energy is NOT conserved**

**by irreversible dissipative non-  
equilibrium processes**

**of  $\text{PTD}(\mathbf{Q}) = 4$**



# *Application Requirements*

1. Must choose the functional form for a 1-form of Action,  $A$ , to define a physical system.
2. **Possibly** constrain the PTD of the Work 1-form,  $W$ , to generate the desired constrained dynamical PDE's.
3. Evaluate the non-equilibrium properties encoded as
  - 3-form of Topological Torsion,  $A \wedge dA$
  - 4-form of Topological Parity,  $dA \wedge dA$ .

# A Universal Example

Use the 4D electromagnetic 1-form of Action:

$$\mathbf{A} = A_k(x,y,z,t)dx^k - \phi(x,y,z,t)dt.$$

Construct the 2-form  $\mathbf{F} = d\mathbf{A} = B_z dx^1 dy^2 \dots - E_z dz^1 dt \dots$

Define:  $\mathbf{E} = -\partial(A_k(\mathbf{x})/\partial t - \text{grad } \phi)$  ;  $\mathbf{B} = \text{curl } A_k$

Construct the Work 1-form:

$$\mathbf{W} = i(\mathbf{V}_4)d\mathbf{A} = -\{\mathbf{E} + \mathbf{V} \times \mathbf{B}\}_k dx^k + \{\mathbf{V} \cdot \mathbf{E}\} dt$$

Then, the coefficients of the 3-form  $d\mathbf{F} = dd\mathbf{A} = 0$ ,  
generate a set of **Maxwell-Faraday** PDE's.

**UNIVERSAL FARADAY INDUCTION** (no **D** or **H**!)

$$\text{Curl } \mathbf{E} + \partial \mathbf{B} / \partial t = \mathbf{0}, \quad \text{div } \mathbf{B} = \mathbf{0}.$$

# A Universal Example

Construct the 3-form  $A^{\wedge}F = A^{\wedge}dA$  (of Topological Torsion)

$$A^{\wedge}dA = i(T_4)(dx^{\wedge}dy^{\wedge}dz^{\wedge}dt)$$

with the 4 component Torsion Direction Field

$$T_4 = - [ \mathbf{E} \times \mathbf{A} + \phi \mathbf{B}, (\mathbf{A} \cdot \mathbf{B}) ].$$

( The 4 divergence of the Torsion vector  $T_4$  is equal to  $-2(\mathbf{E} \cdot \mathbf{B}) = -2\sigma$  .

$$\text{Then, } L_{(T_4)} A = (\mathbf{E} \cdot \mathbf{B}) A = \sigma A$$

$\sigma$  is a Conformality Factor,

and a Homogeneity (including fractals) index of self similarity.

(Suggested as a dissipative extension to Hamilton's principle by RMK in 1974)

# A Universal Example

Construct the 4-form  $F^{\wedge}F = dA^{\wedge}dA$  of Topological Parity:

$$dA^{\wedge}dA = 2(\mathbf{E} \cdot \mathbf{B})(dx^{\wedge}dy^{\wedge}dz^{\wedge}dt).$$

On regions where  $\sigma = (\mathbf{E} \cdot \mathbf{B}) > 0$ ,  
the PTD(A) is 4.

The thermodynamic system is not in equilibrium. And as

$$L_{(T_4)}A^{\wedge}L_{(T_4)}dA = Q^{\wedge}dQ = (-\mathbf{E} \cdot \mathbf{B})^2 A^{\wedge}dA > 0.$$

*The process in the direction of the Torsion vector  $T_4$   
is thermodynamically irreversible*

# *“Entropy” production Rate*

## *Bulk Viscosity*

$$\text{As } L_{(T)} \mathbf{A} = (\mathbf{E} \cdot \mathbf{B}) \mathbf{A} ,$$

$$Q^{\wedge} dQ = (-\mathbf{E} \cdot \mathbf{B})^2 \mathbf{A}^{\wedge} d\mathbf{A} > 0.$$

$$\text{Then } (\mathbf{E} \cdot \mathbf{B})^2 > 0$$

**Entropy Production Rate**

$$= (1/2 \text{ divergence of } \mathbf{T}_4)^2$$

**For Engineers**

**To Reduce Irreversible  
Dissipation and Turbulence**

**For Engineers**

**To Reduce Irreversible  
Dissipation and Turbulence**

**In dynamical systems,  
minimize the co-linearity of  
Translation Acceleration  
and Vorticity**

In Engineering terms

**To Reduce Irreversible  
Dissipation and Turbulence**

In a Plasma Minimize  $\sigma = (\mathbf{E} \cdot \mathbf{B})$



For Engineers

## To Reduce Irreversible Dissipation and Turbulence

In a Navier Stokes fluid Minimize  $(\mathbf{a} \cdot \boldsymbol{\omega})$

$$\frac{\partial \mathbf{V}}{\partial t} + \text{grad } V^2/2 - \mathbf{V} \times \text{curl } \mathbf{V} = - \text{grad} P/\rho + \lambda \text{grad div } \mathbf{V} - \nu \nabla^2 \mathbf{V}$$

Classically,  $\nu$  = shear viscosity, and  $\lambda = (\mu_B - \nu)$  where  $\mu_B$  = Bulk viscosity

Turbulent solutions imply  $\sigma = (\mathbf{a} \cdot \boldsymbol{\omega}) \neq 0!$

# Additional Ideas 1987-2009

## EMERGENCE of Topological Defects

**CONJECTURE:** The thermodynamic Cosmological Astronomical Universe is best represented as a category of very low density radiation plasma of indistinguishable complex wavelets and macroscopic Spinor fields, combined with distinguishable sets of massive particles forming a non-equilibrium system of Pfaff Topological Dimension 4.

**The Thermodynamic PDT4 environment is not a void.**

# Additional Ideas 1987-2009

## EMERGENCE of Topological Defects

**Claim:** The Category theory of Topological Thermodynamics yields examples how continuous **irreversible** dissipative processes can form **emergent** Closed topological defect structures, far from equilibrium, of  $PTD = 3$  (such as **stars**), in an Open thermodynamic environment of  $PTD = 4$ .

**These examples give formal justification to Prologine's conjectures.**

# *Emergence of Stationary States far from equilibrium* by irreversible Processes

**Example:** **Continuous topological evolution** can describe the irreversible evolution on an

**“Open”** symplectic non-equilibrium domain of Pfaff dimension **4**, with evolutionary irreversible orbits being attracted to a contact

**“Closed”** non-equilibrium domain of Pfaff dimension **3**, with an ultimate decay to the

**“Isolated-Equilibrium”** domain of Pfaff dimension **2** or less (integrable Caratheodory surface).

# The Experimental Emergence of Falaco Solitons



Long Lived Topological Defects in a Swimming Pool

# Properties of 3-Form Currents

## J = Processes

$$L_{(J)} \mathbf{A} \Rightarrow \mathbf{Q} \quad L_{(J)} d\mathbf{A} \Rightarrow d\mathbf{Q}$$

### 1. Necessary condition for Reversible Processes

$$\mathbf{Q} \wedge d\mathbf{Q} = \mathbf{0}$$

### 2. Necessary condition for Irreversible Processes

$$\mathbf{Q} \wedge d\mathbf{Q} \neq \mathbf{0}$$

Based on the Pfaff Topological Dimension of  $W$

Thermodynamic **Reversible Processes** imply that the

Heat 1-form,  $Q$  is integrable.  $Q \wedge dQ = 0$

**Extremal:**

Hamiltonian global

$$\text{PTD}(W) = 0, W = 0,$$

**Bernoulli-Casimir:**

Hamiltonian local

$$\text{PTD}(W) = 1, W \text{ exact}$$

**Helmholtz:**

Conservation of Vorticity

$$\text{PTD}(W) = 1, W \text{ closed}$$

Each of these flows are thermodynamically reversible, as

$$dW = 0 = dQ, \text{ implies } Q \wedge dQ = 0.$$

**All of the above processes are Reversible!**  
**These theories make up the bulk of Classical Dynamics!!**

**THESE PROCESSES (classical theories) CAN NOT**  
**faithfully DESCRIBE**  
**TURBULENCE OR**  
**INTERACTION with the ENVIRONMENT**

To be **Irreversible** the process and the thermodynamic system must be non-integrable,  $Q^{\wedge}dQ$  is not zero.

The Topological Torsion (based on  $Q$ )  
is NOT ZERO.  $PTD(Q) > 2$ .



Non-Equilibrium systems are **non-integrable**, in the sense that if there is a solution, it is not unique, or Not Deterministic.

A key artifact of non-equilibrium is the existence of

**Topological Torsion current 3-forms,**  $J_{\text{torsion}} = A \wedge dA,$

**Topological Spin current 3-forms,**  $J_{\text{spin}} = A \wedge G$

**Topological Adjoint current 3-forms,**  $dJ_{\text{adjoint}} = A \wedge J_{\text{spin}}$

# 3-Form Currents = Processes

The Topological Torsion 3-form  $\mathbf{A}^{\mathbf{F}}$  is related to Helicity,

The Topological Spin 3-form  $\mathbf{A}^{\mathbf{G}}$  is related to Spin,

The Adjoint 3-form  $\mathbf{J}_{\text{adjoint}}$  is related to the interaction energy.

All three are related to different species of **dissipative phenomena**, which only occur in non-equilibrium systems.

**The dissipation coefficients  $\sigma$  are equal to the non-zero divergences of the vector coefficients of each 3-form.**

In electromagnetic systems, the 4D dissipation coefficient  $\sigma$  was shown to be equal to  $\mathbf{E} \cdot \mathbf{B}$ ; in hydrodynamics, the 4D dissipation coefficient is called  $\sigma$  "**Bulk viscosity**".

The combination of **Continuous topological evolution**  
and **Parity decay** determine

# A Topological Arrow of Time.

You can describe the **decay** of turbulence **continuously**,  
but NOT the **creation** of turbulence.

# New Ideas 2009-2011

The Not-T0 topologies, without separation axioms, are such that some, if not all, of the subset closures are

**Indistinguishable.**

These topologies appear to describe statistical and quantum-like thermodynamical systems, whose parts, like Bosons and Fermions, are **indistinguishable** -- with equations of “evolution” described by complex diffusion or wave equations with source.

# New Ideas 2009-2011

Stimulated by Cartan's successes for particle systems, the "laws of motion" that describe the continuous evolution of the Not-T0 topologies will be studied by applying the homotopic Lie differential to the exterior differential N-form density,

$$\mu = \rho \{dx^{\wedge}dy^{\wedge}dz^{\wedge}dt\}.$$

The 2<sup>nd</sup> Cartan homotopic evolution formula becomes:

$$L_{(J)} \mu = i(J) d\mu + d(i(J) \mu) = \kappa \mu$$

where  $\kappa$  = the chaotic similarity, or 4D expansion coefficient

# New Ideas 2009-2011

Stimulated by Cartan's successes for particle systems, the "laws of motion" that describe the continuous evolution of the Not-T0 topologies will be studied by applying the homotopic Lie differential to the exterior differential N-form density,

$$\mu = \rho \{dx^{\wedge}dy^{\wedge}dz^{\wedge}dt\}.$$

The 2<sup>nd</sup> Cartan homotopic evolution formula becomes:

$$L_{(\mathbf{J})} \mu = \mathbf{i}(\mathbf{J}) d\mu + d(\mathbf{i}(\mathbf{J}) \mu) = \kappa \mu$$

Leading to a possible statistical (wave) component of entropy, without counting particles.

# New Ideas 2009-2011

Evaluate

$$L_{(\mathbf{J})} \mu = i(\mathbf{J}) d\mu + d(i(\mathbf{J}) \mu) = \kappa \mu$$

to obtain

$$L_{(\mathbf{J})} \mu = 0 + d(i(\mathbf{J}) \mu) = \kappa \mu$$

or

$$\operatorname{div}_4 \mathbf{J} = \kappa - \mathbf{J}^k \partial \ln p / \partial x^k$$

**Cartan's Second Fundamental Equation**

# New Ideas 2009-2011

Solutions to the 2nd fundamental equation can be determined in terms of systems of PDE's that satisfy the fundamental formula:

$$\text{div}_4 \mathbf{J} = \kappa - \mathbf{J}^k \partial \ln p / \partial x^k$$



## **Different PDE Solutions can be recognized as**

**The Wave Equation.**

**The Wave equation with Sources.**

**The Diffusion equation.**

**The Schroedinger equation**

**The Minimal Surface equation.**

**The Dissipation-Interaction equation.**

**The Ginsburg-Landau equation.**

**The Gibbs entropy equation.**

**The Mandelbrot entropy equation for chaos.**

## **Different PDE Solutions can be recognized as**

**The details of such system solutions of PDE's is given in the attached pdf file. A few simple examples are given in that which follows.**

Solutions to the 2<sup>nd</sup> fundamental formula leads to systems of PDE's representing **Not-T0 topologies** of topological indistinguishable subsets.

$$\text{div}_4 \mathbf{J} = \kappa - \mathbf{J}^k \partial \ln p / \partial x^k$$

The term,  $\mathbf{J}^k \partial \ln p / \partial x^k$ , if equal to zero, is a generalization of the **Eikonal Equation**, which can represent propagating discontinuities. (The “discontinuous signals” of the Wave Equation)

# Example 1 The 4D Wave Equation

Consider a set of topological coordinates

$$(\mathbf{x}^k) = [x, y, z, s=ict]$$

with a density 4 form =  $\rho(\mathbf{x}^k) dx^x dy^y dz^z ds$

Identify the density coefficient  $\rho(\mathbf{x}^k)$  with the symbol  $\Psi(\mathbf{x}^k)$  and  
 $\text{grad}_4(\Psi) = \text{grad}_4(\rho)$

Consider the process  $\mathbf{J}$  defined by  $\text{grad}_4(\Psi)$

$$\mathbf{J}(\mathbf{x}^k) = \left( \frac{\partial \Psi}{\partial x}, \frac{\partial \Psi}{\partial y}, \frac{\partial \Psi}{\partial z}, \frac{\partial \Psi}{\partial s} \right)$$

$$\text{div}_4 \mathbf{J}(\mathbf{x}^k) = \frac{\partial^2 \Psi}{\partial x^2} + \frac{\partial^2 \Psi}{\partial y^2} + \frac{\partial^2 \Psi}{\partial z^2} - \frac{\partial^2 \Psi}{\partial s^2}$$

# Example 1 The 4D Wave Equation

**THEN**

$$\text{If } \kappa - \mathbf{J}^k \partial \ln \rho / \partial x^k \Rightarrow 0$$

$$\text{div}_4 \mathbf{J}(x^k) \Rightarrow 0$$

**Therefore**

$$\partial^2 \Psi / \partial x^2 + \partial^2 \Psi / \partial y^2 + \partial^2 \Psi / \partial z^2 - \partial^2 \Psi / c^2 \partial t^2 \Rightarrow 0$$

**THE WAVE EQUATION !**

## Example 2 The Diffusion Equation

Consider a set of topological coordinates

$$(\mathbf{x}^k) = [x, y, z, t]$$

with a density 4 form =  $\rho(\mathbf{x}^k) dx^x dy^y dz^z dt$

Identify the density coefficient  $\rho(\mathbf{x}^k)$  with the symbol  $\Psi(\mathbf{x}^k)$  and  
 $\text{grad}_4(\Psi) = \text{grad}_4(\rho)$

Choose the process current format as:

$$\mathbf{J}(\mathbf{x}^k) = \left( \frac{\partial \Psi}{\partial x}, \frac{\partial \Psi}{\partial y}, \frac{\partial \Psi}{\partial z}, -D \frac{\partial \Psi}{\partial t} \right);$$

$$\text{div}_4 \mathbf{J}(\mathbf{x}^k) = \frac{\partial^2 \Psi}{\partial x^2} + \frac{\partial^2 \Psi}{\partial y^2} + \frac{\partial^2 \Psi}{\partial z^2} - D \frac{\partial \Psi}{\partial t}$$

## Example 2 The Diffusion Equation

**THEN**

$$\text{If } \kappa - \mathbf{J}^k \partial \ln p / \partial \mathbf{x}^k \Rightarrow 0$$

$$\text{div}_4 \mathbf{J}(\mathbf{x}^k) \Rightarrow 0$$

**Therefore**

$$\partial^2 \Psi / \partial \mathbf{x}^2 + \partial^2 \Psi / \partial \mathbf{y}^2 + \partial^2 \Psi / \partial \mathbf{z}^2 - \mathbf{D} \partial \Psi / \partial \mathbf{t} \Rightarrow 0$$

**THE DIFFUSION EQUATION !**

### Example 3 The Schroedinger Equation

Consider a set of topological coordinates

$$(\mathbf{x}^k) = [x, y, z, t]$$

with a density 4 form =  $\rho(\mathbf{x}^k) dx^x dy^y dz^z dt$

Choose the process current as:

$$\mathbf{J}(\mathbf{x}^k) = (\partial\Psi/\partial x, \partial\Psi/\partial y, \partial\Psi/\partial z, -h/i \Psi);$$

$$\text{div}_4 \mathbf{J}(\mathbf{x}^k) = \partial^2\Psi/\partial x^2 + \partial^2\Psi/\partial y^2 + \partial^2\Psi/\partial z^2 - h/i \partial\Psi/\partial t$$



## Example 3 The Schroedinger Equation

**THEN**

$$\text{If } \kappa - \mathbf{J}^k \partial \ln \rho / \partial x^k \Rightarrow v\Psi$$

$$\text{div}_4 \mathbf{J}(\mathbf{x}^k) \Rightarrow v\Psi$$

**Therefore**

$$\partial^2 \Psi / \partial x^2 + \partial^2 \Psi / \partial y^2 + \partial^2 \Psi / \partial z^2 - \hbar / i \partial \Psi / \partial t = v\Psi$$

**THE SCHROEDINGER EQUATION !**

## Summary:

THE CATEGORY THEORY OF **T0** TOPOLOGY (PARTICLES) AND **NON-T0** TOPOLOGY (WAVES), and THE POSSIBLE INTERACTIONS OF THE TWO TOPOLOGICAL TYPES, HAS JUST BEGUN.

The two **universal** equations of Topological Evolution are:

$$L(\mathbf{J}) \mathbf{A} = \mathbf{i}(\mathbf{J})d\mathbf{A} + d(\mathbf{i}(\mathbf{J})\mathbf{A}) = \mathbf{W} + d(\mathbf{U}) = \mathbf{Q}$$

$$\text{div}_4 \mathbf{J} = \mathbf{\kappa} - \mathbf{J}^k \partial \ln \rho / \partial x^k$$

I AM CONVICTED THAT THE BOSE – EINSTEIN  
AND THE FERMI DIRAC DISTRIBUTIONS CAN BE  
DEDUCED FROM

$$\text{div}_4 \mathbf{J} = \kappa - \mathbf{J}^k \partial \ln p / \partial x^k$$

CAN YOU HELP ME ???

# Thanks for your interest

Contact Professor R. M. Kiehn at

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# Non-Equilibrium **Systems** and Irreversible **Processes**

Paperback Monographs @ <http://www.lulu.com/kiehn> Free PDF files: [toptorsion@aol.com](mailto:toptorsion@aol.com)

Vol 1

## Non-Equilibrium Thermodynamics

Sliding Ball => Rolling Ball



Irreversible Continuum: Topological Evolution of Pfaff Topological Dimension  $2n-2$  to long lived states for Non-equilibrium and of Pfaff dimension  $2n+1$ .

R. M. Kiehn

Vol 2

## Falaco Solitons Cosmology and the Arrow of time



Photo courtesy David Radabaugh

R. M. Kiehn

Vol 3

## Wakes Coherent Structures and Turbulence



Photo Courtesy Paul Rowan [www.airbase.net](http://www.airbase.net)

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Vol 4

## Plasmas and Non-Equilibrium Electrodynamics



Long lived ionized plasma ring in the aftermath of a nuclear explosion.

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Vol 5

## Topological Torsion and Macroscopic Spinors

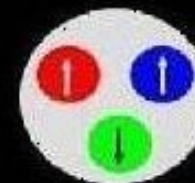


Helicity of the 5-D and  
Macroscopic Spinors

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Vol 6

## The Universal Effectiveness of Topological Thermodynamics



Associativity Preorder in Quarks

R. M. Kiehn

From the perspective of Continuous Topological Evolution

More to come

# **50 Significant Results**

**based on the**

**The Category of  
Topological Thermodynamics**



# Significant Results

- The method of Pfaff Topological Dimension explains topological differences between equilibrium  $PTD < 3$  (connected) and non-equilibrium  $PTD > 2$  (disconnected) thermodynamic systems of distinguishable “particles”.
- The method explains the topological distinctions between thermodynamically reversible integrable ( $Q^{\wedge}dQ=0$ ) and irreversible non-integrable ( $Q^{\wedge}dQ \neq 0$ ) processes.
- The method explains topological distinctions between thermodynamic systems based on **distinguishable** sets (T0 topologies) and thermodynamic systems based upon statistical distributions of **indistinguishable** sets (Not T0 topologies).

# Significant Results

- The method explains that **Entropy** can consist of two simultaneous components: one component is based on **distinguishable** sets (massive particles) and the other component is based upon **indistinguishable** sets (radiation).
- The method demonstrates that the First Law of Thermodynamics is a statement of cohomology, where the difference of the two, not necessarily exact, differential 1-forms of Heat **Q** and Work **W** are topologically equal to an exact differential of internal energy, **dU**.
- The method demonstrates that the **decay** of turbulence can be described in terms of a topologically continuous process, but the **creation** of turbulence can not.

# Significant Results

- The method indicates that the top down nested sequence of **Open**, **Closed**, **Isolated-Equilibrium** and **Equilibrium** thermodynamic domains is in 1-1 correspondence with the Pfaff Topological dimension,  $[4,3,2,1]$  generated by any 1-form of Action **A** used to create the Cartan topological structure.
- The environment of the universe is not a vacuum (a void empty set  $PTD(\mathbf{A}) = 0$ ), but is best described as the disconnected **Open** system of  $PTD(\mathbf{A}) = 4$  in which sub-domains of  $PTD < 4$  are distinguishable topological defects.
- The **Equilibrium** systems are of  $PTD(\mathbf{A}) = 1$ . The **Isolated Equilibrium** systems are of  $PTD(\mathbf{A}) = 2$ . The **Closed** systems are of  $PTD(\mathbf{A}) = 3$ .

# Significant Results

- The Pfaff Topological Dimension is equal to the number of non-zero entries in the Pfaff Sequence:

$$\{A, dA, A \wedge dA, dA \wedge dA\}$$

of the 1-form used to create the Cartan topological structure.

- Non-Equilibrium **Open** systems of Pfaff Topological Dimension 4 (PTD=4) can exhibit local decay (in finite time) to Closed non-equilibrium systems of PTD =3, thereby justifying Prigogine's conjecture of emergence to states far from equilibrium by means of dissipative processes. These metastable PTD=3 **Closed** states are topological defects in the Open systems of PTD=4.

# Significant Results

- Topological change is a necessary condition for thermodynamic irreversibility.
- The combination of continuous, non-homeomorphic, processes of topological evolution establish a logical basis for thermodynamic irreversibility and the arrow of time.
- Any synergetic system of parts defines a topology such that the **Category of Topological Thermodynamics** is applicable to economic systems, political systems, as well as to biological systems. Such systems admit long-lived states far from equilibrium (adulthood). However, the theory can be used to design subtle perturbations to such systems such that their metastable lifetimes can be extended or destroyed.

# Significant Results

- **C2 Continuous Topological Evolution permits irreversible processes,  $Q^dQ \neq 0$ . Segmented C1 processes, which approximate C2 processes, can be reversible,  $Q^dQ = 0$ .**
  - **On odd-dimensional spaces, sequential C1 (translational) processes can be thermodynamically reversible, while intransitive C2 processes (rotation and expansion with a fixed point) can be thermodynamically irreversible.**
  - **The relativistic Twin Paradox is resolved if the evolutionary paths of each twin are defined by processes that cause (unequal) topological change. If there is no topological change, or if each twin suffers the same topological change, there is no disparate biological aging.**

# Significant Results

- **Adiabatic processes are transverse to the heat 1-form,  $(i(\rho V_4)Q) = 0$ . Adiabatic processes need not be quasi-static, and can be reversible or irreversible. Topological Transitions can occur in finite (often short) time.**
  - **The work 1-form,  $W$ , is always transverse to the 4D process direction field,  $\rho V_4$ , but the heat 1-form,  $Q$ , may or may not be transverse, The heat 1-form,  $Q$  can have longitudinal components in the direction of the process, corresponding to irreversible dissipation.**
  - **Engineers should be guided by the universal concept of minimizing longitudinal heat in order to improve efficiency.**

# Significant Results

- For fluids, this idea can be translated to reducing vorticity components in the direction of fluid accelerations (by using wing tip tabs on commercial jets).
- For non-equilibrium systems, the 3-form of Topological Torsion is not zero:  $\mathbf{A} \wedge d\mathbf{A} = i(T_4) dx \wedge dy \wedge dz \wedge dt \neq 0$ .  
 $T_4$ , is deduced intrinsically from the 1-form,  $\mathbf{A}$ , that encodes the thermodynamic system, and can be used as special process current density,  $\rho T_4$ , that defines an **irreversible** process.
- For  $PTD=3$  "closed" thermodynamic systems, the process current has zero divergence, and the 4D volume element is a conformal invariant (any  $\rho$ ). This result is the space-time extension of the Liouville theorem that preserves the phase-space volume element in classical theory.



# Significant Results

- For a  $PTD=4$  "open" thermodynamic systems, the Topological Torsion vector does NOT have zero divergence, and so the process current  $\rho T_4$  may not be closed for arbitrary  $\rho$  (that is, the divergence of the process current is not zero).
  - The combination of continuity and non-equilibrium,  $PTD(Q)>2$ , requires a causal direction to the arrow or time.
  - $PTD(A)=3$  domains can have a local basis in terms of one complex **Spinor** pair with complex conjugate eigenvalues, and one real vector with eigenvalue zero. The eigenvalue 0 state can represent metastable (long-lived) configurations far from equilibrium. Such domains are locally "contact" manifolds.

# Significant Results

- A key artifact of non-equilibrium is the existence of **Topological Torsion** current 3-forms, **Topological Spin** current 3-forms and **Topological Adjoint** current 3-forms.

All of these current 3-forms are similar to the Amperian charge current 3-form of electromagnetic theory, but are related to different species of dissipative phenomena, and occur only in non-equilibrium systems.

Dissipation coefficients are related to the non-zero divergences of the vector coefficients of the various 3-forms. For example, in electromagnetic systems, the dissipation coefficient is proportional to  $\mathbf{E} \circ \mathbf{B}$  (see Vol 4) ; in hydrodynamics, the dissipation coefficient is called "**Bulk Viscosity**" (see Vol 3).

# Significant Results

- $\text{PTD}(\mathbf{A})=4$  domains can have a local basis in terms of two complex **Spinor** pairs. Locally, such domains are symplectic manifolds. In such symplectic domains there exist local density distributions (subdomains),  $\rho$ , such that the divergence of any process current is zero in that subdomain.

Such subdomains are metastable contact manifolds. In other words, the symplectic domain can contain defect structures in the form of contact subdomains. It can be demonstrated in terms of continuous topological evolution that such local density distributions, which define a "stationary" state, can **emerge** as a topological defect in a  $\text{PTD}=4$  system, by means of a dissipative processes.

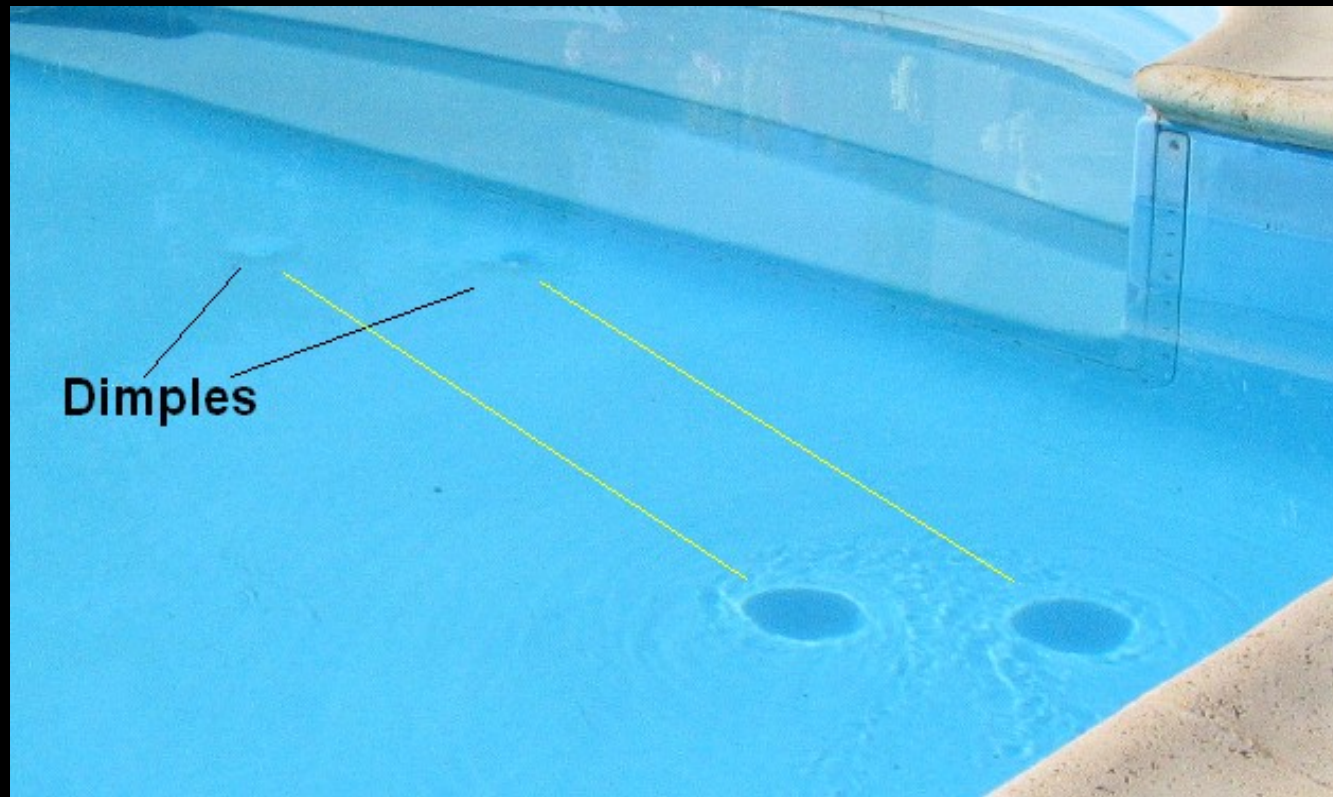
# Emergence of Topological Defects



Long Lived Topological Defects in a Swimming Pool

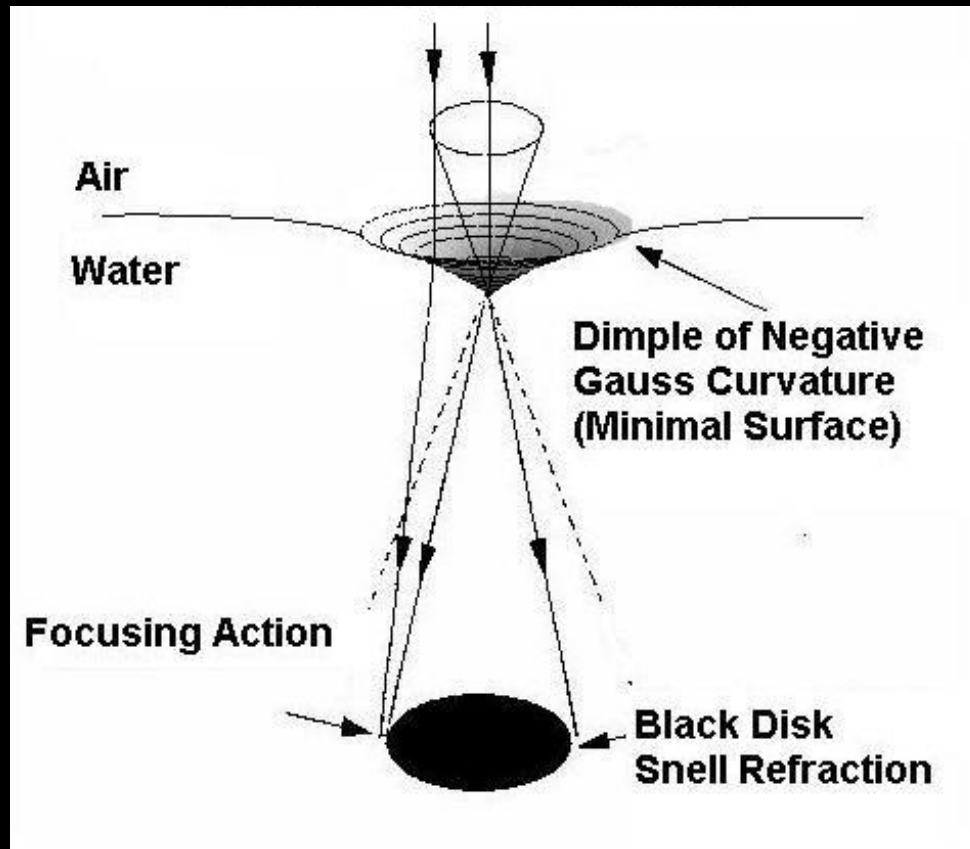
# Emergence of Topological Defects

FALACO SOLITONS Movie by D. Radabaugh



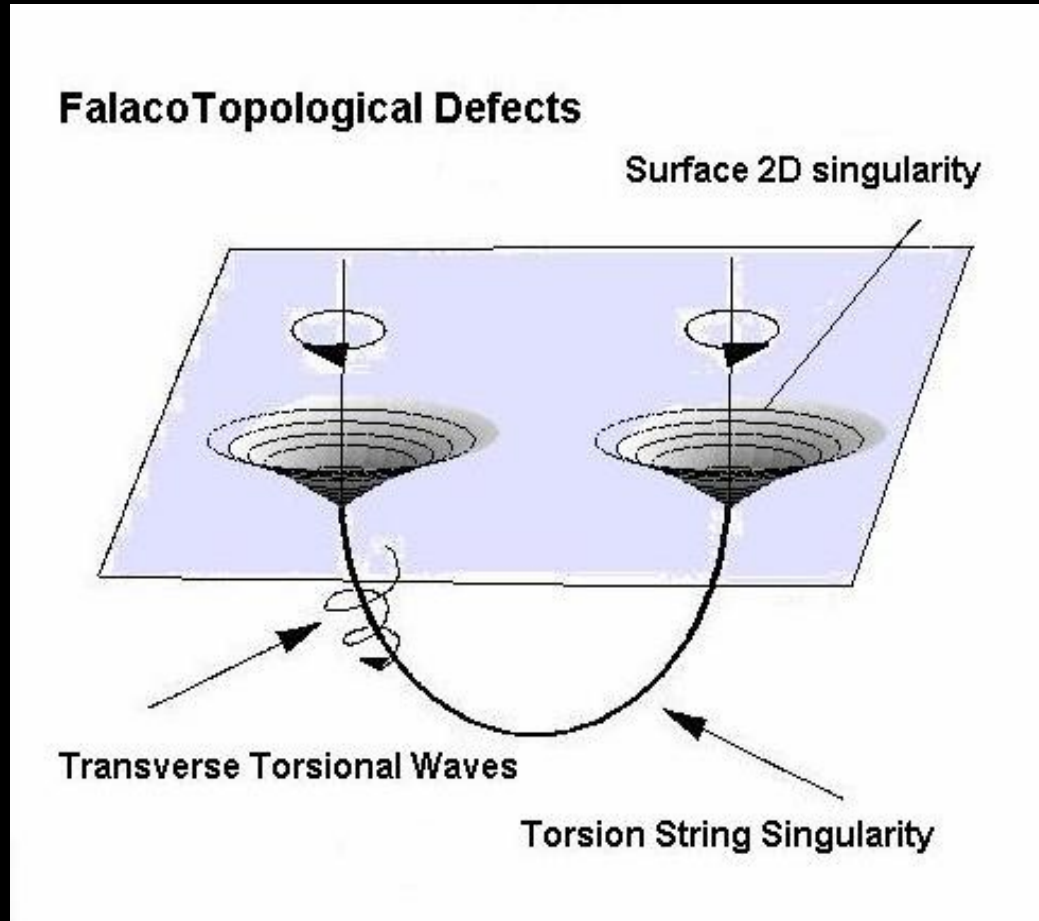
**Solar Elevation about 30 degrees (See movie at  
<http://www22.pair.com/csdc/download/spotsmovie.avi>)**

# Emergence of Topological Defects



Snell refraction of Falaco Soliton Spin Pairs

# Emergence of Topological Defects



The first measurable Torsion String coupling between branes

This real world effect has been ignored by string theorists !!!

# Significant Results

The 3-form of Adjoint Current,  $\mathbf{J}_{\text{adj}}$ , like the Torsion Current, can be constructed entirely from the topological features of the thermodynamic system as determined by the functional coefficients of the 1-form of Action,  $\mathbf{A}$ .  $\mathbf{J}_{\text{adj}}$  also admits an infinity of integrating factors. However, when the 1-form coefficients,  $A_k$  are divided by a Holder Norm,  $\lambda$ , of homogeneity index one (known as the Gauss map), then the vector components of  $\mathbf{J}_{\text{adj}}$  are defined as  $|\mathbf{J}_{\text{adj}}\rangle = [\text{Jacobian}(A_k/\lambda)]^{\text{adj}} \circ |A_k/\lambda\rangle$ . The 3-form is defined in terms of vector components as  $\mathbf{J}_{\text{adj}} = i(\mathbf{J}_{\text{adj}}\}) dx^{\wedge} dy^{\wedge} dz^{\wedge} dt$ . It has zero divergence globally, and a pre-image 2-form  $\mathbf{G}_{\text{adj}}$ , such that  $d\mathbf{G}_{\text{adj}} = \mathbf{J}_{\text{adj}}$ . It can be demonstrated that  $\mathbf{J}_{\text{adj}}$  is related to the cubic curvature of the shape matrix, leading to the idea that the source of electromagnetic charge is related to cubic curvatures, similar to the idea that mass is related to quadratic curvatures.



# Significant Results

- Examples of thermodynamic systems can be given to demonstrate that the conjectured format of the London Current of superconductivity, where  $\mathbf{J} = \chi \mathbf{A}$ , can be deduced as an emergent consequence of the Topological Theory of Thermodynamics (see Vol 5).
  - Examples can generate a Spin Current 3-form,  $\mathbf{S} = \mathbf{A} \wedge \mathbf{G}$ , where, formally, the Spin current is proportional to the Lorentz force (the space-time components of the Work 1-form,  $W$ ). This is a new interpretation of an old result,  $\mathbf{J} = \sigma(\mathbf{E} + \mathbf{V} \times \mathbf{B})$ , which is Ohm's law. The new part is due to the idea that the dissipation is due to Spin Currents and the transport of collective spins,  $\mathbf{A} \wedge \mathbf{G}$ .

# Significant Results

- The topological structure of domains of Pfaff dimension 2 or less creates a connected, but not necessarily simply connected topology. Evolutionary **predictive solution uniqueness** is possible.
- The topological structure of domains of Pfaff dimension 3, or more, creates a disconnected topology of multiple components. If solutions to a particular evolutionary problem exist, then the **solutions are not unique**. Envelope solutions, such as Huygen wavelets and propagating tangential discontinuities (called signals, or wakes) are classic examples of solution non-uniqueness. **Topological Torsion is an artifact of non-uniqueness, and must be non-zero if a hydrodynamic system exhibits turbulence.**

# Significant Results

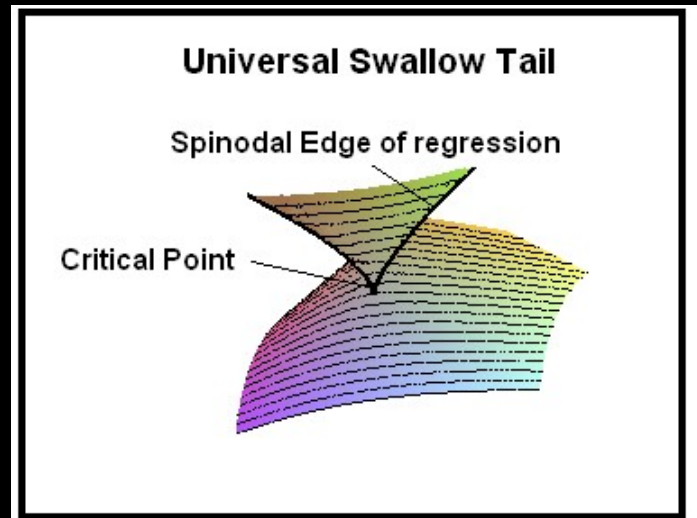
- All Hamiltonian, Symplectic-Bernoulli and Helmholtz processes **are thermodynamically reversible**. In particular, the work 1-form, **W**, created by Hamiltonian processes is of Pfaff Topological Dimension 1 or less. In all reversible cases the Work 1-form is closed,  $d\mathbf{W}=0$ . **The evolutionary equations in such cases are time reversal invariant.**
- The assumption of uniqueness of evolutionary solutions (which implies the Pfaff Topological Dimension of the thermodynamic system be equal to 2 or less), or the assumption of homeomorphic evolution, have imposed constraints upon classical mechanics that eliminate any time asymmetry, and preclude the existence of isotropic **macroscopic spinors**.

# Significant Results

- The Lie differential acting on differential forms is not necessarily the same as a linear affine covariant differential acting on differential forms. The covariant differential always defines an adiabatic process, where the Lie differential does not.
- The particle view of thermodynamics is based upon the 1-form of action whose coefficients,  $[A_m]$ , admit a **correlation** Jacobian matrix  $[\partial A_m / \partial x^n]$ .
- The statistical point of view is based upon a current N-1 form whose coefficients,  $[C^m]$ , admit a collineation Jacobian matrix,  $[\partial C^m / \partial x^n]$ , with a trace = to a divergence of  $[C^m]$ .

# Significant Results

- On spaces of  $PTD=4$ , the Correlation Jacobian matrix has a characteristic polynomial that defines an quartic equation of state in terms of Cayley-Hamilton similarity invariants. The characteristic polynomial produces an implicit hypersurface function that can have envelopes and edges of regression in the format of the Gibbs function for a (universal and deformable) **van der Waals gas**.

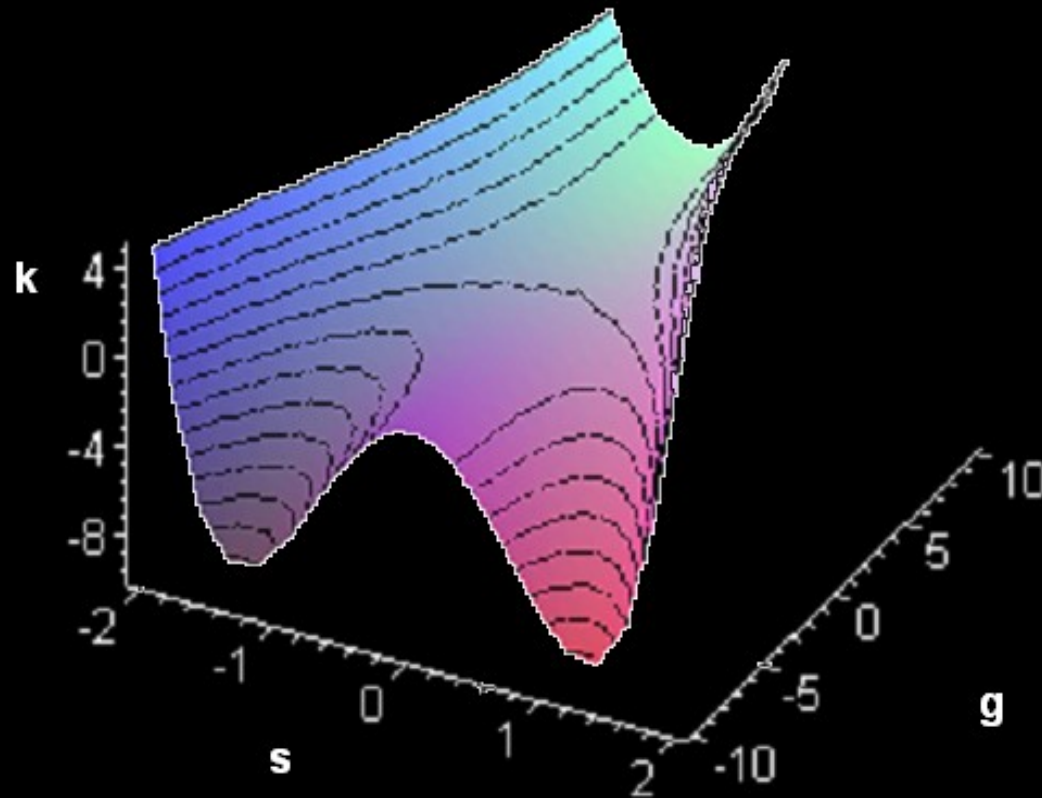


# Significant Results

- The 4D Correlation Jacobian matrix can be mapped into a reduced characteristic polynomial representing a quartic Universal Thermodynamic Phase function, with an envelope, which, below the critical point, has the features of a Higgs potential, made famous by string theory.

The reduced thermodynamic phase function is a hypersurface with Zero mean curvature (a minimal surface) and many properties of a **van der Waals gas**.

# Universal Topological Thermodynamic Phase Function



A van der Waals gas with a **Higgs potential**,  
An Envelope of a 4D Cayley-Hamilton characteristic polynomial

# Significant Results

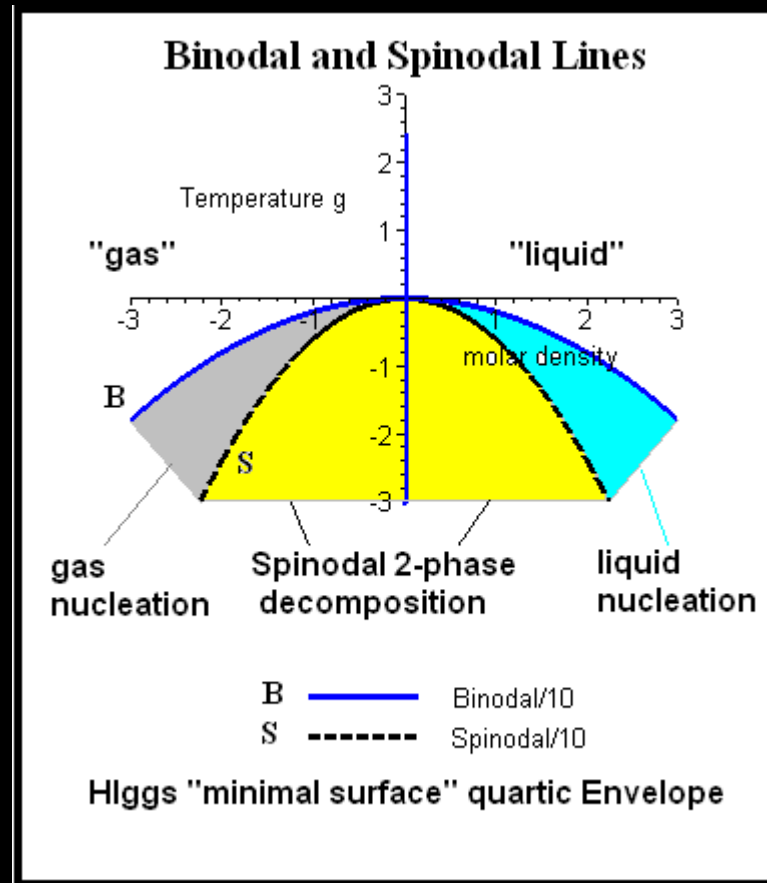
- The reduced phase function yields analytic expressions for the critical point and the binodal and spinodal lines in terms of the similarity invariants of the correlation matrix.

The same technique can be used to determine “phase transitions” and possible bifurcations in dynamical systems.



# Significant Results

- The reduced phase function yields analytic expressions for the critical point and the binodal and spinodal lines in terms of the similarity invariants of the correlation matrix.



# Significant Results

- The collineation matrix can have complex eigenvalues even though the maximal rank matrix elements are real. These complex eigenvectors have a zero norm, and are generators of conjugate minimal surfaces in 4D, that have both statistical and fractal components. Cartan defined such sets as Isotropic Spinors (Majorana, not Dirac) Spinors. Pairs of non-colinear Spinors define an area, but the norm of each Spinor is zero!

The hypersurface minimal surface can be generated by a holomorphic function that includes both the Gibbs entropy and a Mandelbrot fractal germ,  $\theta = (z \ln z - z) + (a \pm bz^2)$ . The third partial derivative leads to conjugate pairs of minimal surfaces, and the Mandelbrot term vanishes. All functional iterates remain holomorphic and generate minimal surfaces with fractal boundaries.

# Significant Results

- For a **PTD=4** "open" thermodynamic systems, the Topological Torsion vector does not have zero divergence,

$$\mathbf{A}^{\wedge}d\mathbf{A}=\mathbf{i}(\mathbf{T}_4)d\mathbf{x}^{\wedge}d\mathbf{y}^{\wedge}d\mathbf{z}^{\wedge}d\mathbf{t} \neq \mathbf{0}.$$

but the process current  $\rho\mathbf{T}_4$  may not be closed for arbitrary  $\rho$  (that is, the divergence of the process current is not zero).

$$d(\mathbf{A}^{\wedge}d\mathbf{A}) \neq \mathbf{0}$$

This result is an extension of the Vlasov equation; the 4D differential volume element is expanding or contracting. Such processes in the direction of  $\mathbf{T}_4$  are irreversible and dissipative.

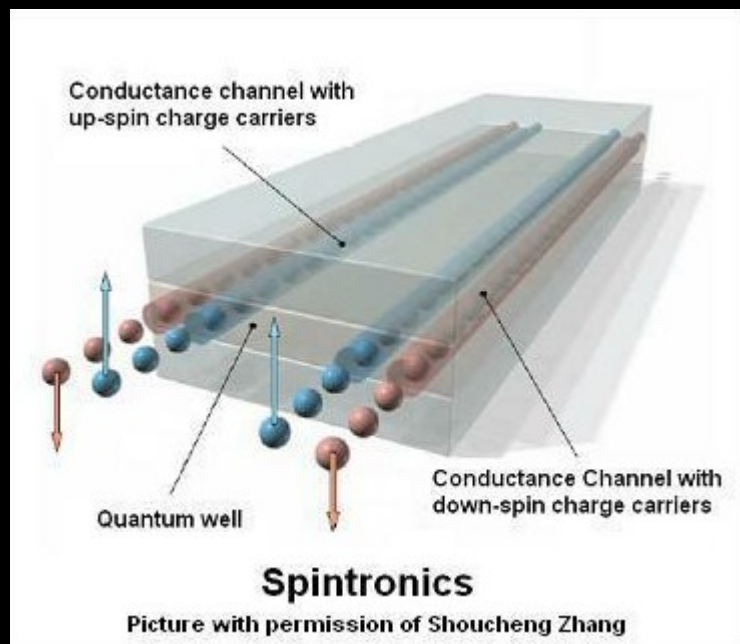
There exist many integrating factors that will produce zero divergence of the 3 form, such that  $d(\rho\mathbf{A}^{\wedge}d\mathbf{A})=0$ .

# Significant Results

- **Topological fluctuations** can be induced by processes that have components in terms of thermodynamic macroscopic Spinors. Thermodynamic **Macroscopic Spinors** are non-zero complex eigenvectors with complex eigenvalues (and zero quadratic form) of the antisymmetric 2-form (or matrix) representing the "Limit Points",  $dA$ , of the 1-form of Action,  $A$ . Such Macroscopic Spinors are capable of representing minimal surface conjugate pairs.

# Significant Results

- **Topological Insulators** correspond to Impedances defined in terms of quantized Topological Spin, and Spin pairs coupled via Falaco Solitons. If the divergence of the Spin 4vector vanishes, the **Spin Current** is time reversal invariant.



# Significant Results

- The fact that any synergetic system of parts in effect defines a topology implies that the universal method of topological thermodynamics is **applicable to economic systems, political systems, as well as to biological systems**. Such systems admit long-lived states far from equilibrium (adulthood). However, the theory can be used to design subtle perturbations to such systems such that their metastable lifetimes can be extended or destroyed.

# Significant Results

- The thermodynamic processes that lead to self-similarity of a Current 3-form  $L_{(j)}C = \sigma C$  can generate **fractals** and **holographic** effects where small neighborhoods replicate the whole, approximately. The homogeneity coefficient is the trace of the Jacobian Collineation:  $\sigma = \text{Trace}[\partial C^m / \partial x^n]$ , or the divergence of the Process vector field.

# Significant Results

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- A turbulent thermodynamic **cosmology** can be constructed in terms of a dilute **non-equilibrium van der Waals gas** near its critical point.



## Cosmology as a non-equilibrium Van der Waals Gas explains

- a.) The granularity of the night sky as exhibited by stars and galaxies due to density fluctuations near the critical point, and the Newtonian law of gravitational attraction proportional to  $1/r^2$  as a correlation between fluctuations (due to Lev Landau ).

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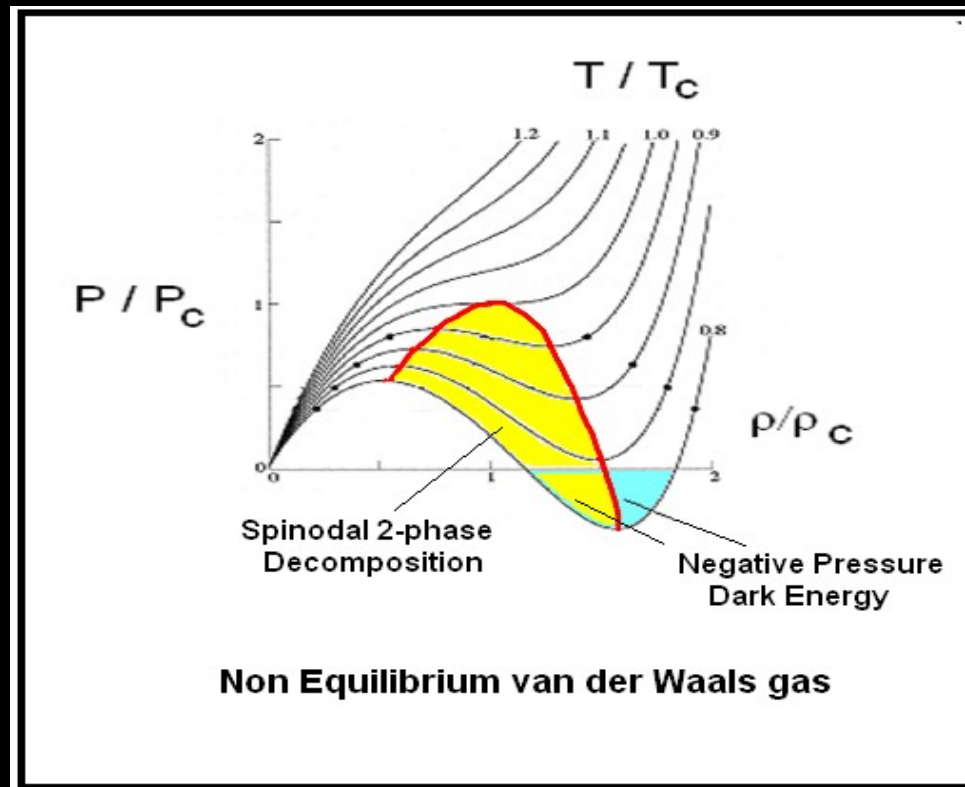
b.) The conformal expansion of the universe is an irreversible phenomenon-associated with Quartic similarity invariants in the thermodynamic phase function, and dissipation related to conformal deformations of the 4D volume element.

## Cosmology as a non-equilibrium Van der Waals Gas explains

- c.) The possibility of domains of negative pressure (explaining what has recently been called "**dark energy**") are due to a classical "Higgs" mechanism for aggregates below the critical temperature.

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## Cosmology as a non-equilibrium Van der Waals Gas explains

- d.) The possibility of domains of negative **temperature** (explaining what has recently been called "**dark matter**") are due to macroscopic collective states of ordered spins. The conjecture is that **Positive temperature radiates, Negative temperature does not.**

Conjecture: **Black holes could be negative temperature states of collective spins.**

## **Cosmology as a non-equilibrium Van der Waals Gas explains**

- e.) The possibility of domains where gravitational effects (quadratic similarity invariants, or 2nd order Gauss curvature effects) appear to be related to entropy and temperature properties of the thermodynamic system, and where **cubic curvature effects due to Spin and Adjoint current 3-forms could impede gravitational collapse.**

## Cosmology as a non-equilibrium Van der Waals Gas explains

- e.) The possibility of domains where gravitational effects (quadratic similarity invariants, or 2nd order Gauss curvature effects) appear to be related to entropy and temperature properties of the thermodynamic system, and where **cubic curvature effects due to Spin and Adjoint current 3-forms could impede gravitational collapse.**
  
- f.) Black Holes (generated by Petrov Type D solutions in gravitational theory) are to be related to **Minimal Surface** solutions to the Universal thermodynamic 4th order Phase function.

# Thanks for your interest

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Vol 1

## Non-Equilibrium Thermodynamics

Sliding Ball => Rolling Ball



Irreversible Continuum: Topological Evolution of Pfaff Topological Dimension  $1n-2$  to long lived states for Non-equilibrium and of Pfaff dimension  $2n+1$ .

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Vol 2

## Falaco Solitons Cosmology and the Arrow of time



Photo courtesy David Radabaugh

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Vol 3

## Wakes Coherent Structures and Turbulence



Photo Courtesy Paul Rowan [www.aircraft.net](http://www.aircraft.net)

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Vol 4

## Plasmas and Non-Equilibrium Electrodynamics



Long lived ionized plasma ring in the aftermath of a nuclear explosion.

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Vol 5

## Topological Torsion and Macroscopic Spinors

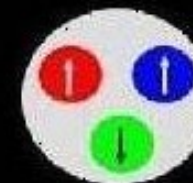


Helicity of the 3-D and  
Macroscopic Spinors

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Vol 6

## The Universal Effectiveness of Topological Thermodynamics



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From the perspective of Continuous Topological Evolution

