From the point of view of Continuous Topological Evolution

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This is an automated 20 minute slide presentation

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My presentation is somewhat low-key to an audience of topologists, but remember, I interact with engineers and scientists that have a limited (if any) training in topology.

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- Almost all of my engineering audience has experience with thermodynamics. I find it to be the universal theory.
- Long ago (1964) I rejected any attempts to describe the dynamics of irreversible processes in terms of equivalence classes of geometric diffeomorphisms (tensors).
- Early on, I concluded that **topological change** is a necessary condition for thermodynamic irreversibility.

It now appears that the topological perspective of thermodynamics gives universal insight into many non-equilibrium concepts associated with the emergence of metastable states, digital topology, plasmas and turbulent flows, biology, chemistry, metallurgy, fuzzy logic, holography and even the cognitive sciences.

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As I told my students, there are two motivations and opportunities for such a discipline:



WNP = Win Nobel Prize

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EBB = Earn Big Bucks (\$)

A Thermodynamic system can be encoded in terms of an exterior differentiable 1-form of Action, A, per unit "mole".

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A Thermodynamic Process can encoded by an ordered array (a vector) of functions that form the coefficients, J, of an N-1=3-form Current.

Topological Thermodynamics is defined in terms of **exterior differential forms** evaluated on ordered classes of differential varieties {x,y,z,t; dx,dy,dz,dt}.

Topological Thermodynamics is defined in terms of **exterior differential forms** evaluated on ordered classes of **differential varieties** {x,y,z,t; dx,dy,dz,dt}.

The ordered class is defined in terms of C1 maps,

from $\{x^k, d x^k\}$ to $\{y^k, d y^k\}$

These maps are not diffeomorphisms, and do not require the geometric constraints of an inverse,

for either ϕ , or d ϕ .

The formal Topological structure of a universal theory of Thermodynamics based on Exterior Differential forms is a

Kolmogorov T₀ Topology

Formally, this topology is quite interesting for many demonstrable reasons. All of the singletons of the topology are not closed. Warning: the Kolmogorov topology is NOT a metric topology, NOT a Hausdorff topology, and even does NOT satisfy the separation axioms that define a T₁ topology

Kolmogorov Topology T₀ of 4 points

```
Table 1. The CT4 Topology of 4 points
                                       X = \{a, b, c, d\}
                        Basis subsets \{a\}, \{a, b\}, \{c\}, \{c, d\}
     CT4\{open\}: \varnothing, \{a\}, \{c\}, \{a, b\}, \{c, d\}, \{a, c\}, \{a, b, c\}, \{a, c, d\}, X
    CT4\{closed\}: X, \{b, c, d\}, \{a, b, d\}, \{c, d\}, \{a, b\}, \{b, d\}, \{d\}, \{b\}, \varnothing
 Subset S
                   Limit Pts
                                     Interior-Exterior
                                                                  Boundary
     ø*
                                                        X
                        \{b\}
                                       \{a\}
                                                     \{c,d\}
                                                                        \{b\}
      \{a\}
                                                                                       \{a,b\}
      \{b\}
\{c\}
                                                                       \{b\}
                                                    \{a, c, d\}
                                                                                        {b}
                                       \{c\}
                                                                       \{d\}
                        \{d\}
                                                      \{a,b\}
                                                                                       \{c,d\}
   d
a, b
                                                     \{a, b, c\}
                                                                       \{d\}
                                                                                        \{d\}
                        {b}
                                        \{a,b\}
                                                                                      \{a,b\}
                                                     \{c,d\}
                     \{b\}, \{d\}
                                                                                         X
     \{a, c\}
                                         \{a,c\}
                                                                      \{b, d\}
                                                      \{c\}
                                        \{a\}
                                                                      \{b, d\}
                                                                                     \{a, b, d\}
     a, d
                        \{b\}
                                                      \{a\}
                                         \{c\}
                                                                                     \{b, c, d\}
    \{b,c\}
                                                                      \{b,d\}
    b, d
                                                     \{a,c\}
                                                                      \{b, d\}
                                                                                       \{b,d\}
    \{c, d\}^*
                        \{d\}
                                        \{c,d\}
                                                      \{a, b\}
                                                                                       \{c,d\}
    a, b, c
                     \{b\}, \{d\}
                                     \{a, b, c\}
                                                                       {b}
   \{b, c, d\}
                        \{d\}
                                      \{c,d\}
                                                      \{a\}
                                                                                     \{b, c, d\}
                                                                        \{b\}
\{d\}
   \{a, c, d\}
                     \{b\}, \{d\}
                                     \{a, c, d\}
                                                                                        X
   \{a, b, d\}
                                      \{a,b\}
                                                       \{c\}
                                                                                     \{b, c, d\}
                        {b}
\{a, b, c, d\}
                     \{b\}, \{d\}
                                                                                         X
```

Kolmogorov Topology dual T*₀ of 4 points

```
Table 2. A DUAL CT4* Topology of 4 points
                                     X = \{a, b, c, d\}
                       Basis subsets \{b\}, \{a, b\}, \{d\}, \{c, d\}
 \text{Dual } CT4\{open\}: X, \{b,c,d\}, \{a,b,d\}, \{c,d\}, \{a,b\}, \{b,d\}, \{d\}, \{b\}, \varnothing \\
Dual CT4\{closed\}: \emptyset, \{a\}, \{c\}, \{a, b\}, \{c, d\}, \{a, c\}, \{a, b, c\}, \{a, c, d\}, X
  Subset S
                   Limit Pts Interior-Exterior Boundary
                                                                                   Closure
                                                       X
                                                                     \{a\}
      \{a\}
                                                                                      \{a\}
                                                 \{b, c, d\}
                        \{a\}
                                        \{b\}
                                                  \{c,d\}
                                                                     \{a\}
                                                                                    \{a,b\}
                                                                     \{c\}
                                                 \{a, b, d\}
                                                                                      \{c\}
                        \{c\}
                                        \{d\}
                                                   \{a,b\}
                                                                                     \{c, d\}
    \{a,b\}^*
                                       \{a,b\}
                                                    \{c,d\}
                                                                                     \{a, b\}
     \{a,c\}
                                                   \{b,d\}
                                                                    \{a,c\}
                                                                                    \{a,c\}
                         \{c\}
     a, d
                                        \{d\}
                                                                    \{a,c\}
                                                                                   \{a, c, d\}
                         \{a\}
                                                    \{d\}
                                        \{b\}
     \{b,c\}
                                                                    \{a,c\}
                                                                                   \{a, b, c\}
     \{b,d\}
                     \{a\}, \{c\}
                                        \{b, d\}
                                                                    \{a, c\}
                                                                                       X
    \{c, d\}^*
                                                    \{a,b\}
                                                                                   \{c, d\}
                                        \{c, d\}
                         \{c\}
     [a, b, c]
                         \{a\}
                                        \{a,b\}
                                                      \{d\}
                                                                     \{c\}
                                                                                   \{a, b, c\}
                                                                     \{a\}
                     \{a\}, \{c\}
                                                                                       X
    \{b, c, d\}
                                       \{b, c, d\}
                                                      {b}
    \{a, c, d\}
                         \{c\}
                                        \{c,d\}
                                                                     \{a\}
                                                                                   \{a, c, d\}
   \{a, b, d\}
                      \{a\}, \{c\}
                                                                                       X
                                       \{a, b, d\}
                                                                                       X
  \{a, b, c, d\}
                                           X
```

The poset of the Kolmogorov topologies, T₀ and T*₀,

$$R = T_0 X T_0^*$$

creates the

Discrete Alexandroff T₁ topology of 4 points.

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To and T* are the topologies of Continuous Fields

T1 is the topology of Discrete Quanta (BITS)

The DISCRETE Alexandroff T1 topology can be partitioned into two CONTINUOUS Kolmogorov topologies, T0 and T*0.

$$\{T1\} => \{T0 + T*0\}$$

{Partitioned Particles} => {Interaction Fields}

Continuous Topological Evolution

Let the Topology of initial state be T1
Let the Topology of the final state be T2

Continuous Topological Evolution

Let the Topology of initial state be T1
Let the Topology of the final state be T2

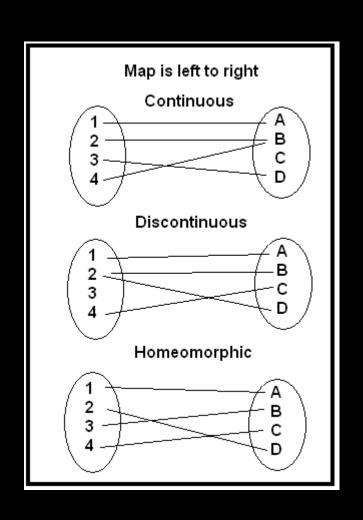
Topological change is continuous iff for the map φ:T1⇒T2
the Limit points of T1 are included in the Closure of T2.

Continuous Topological Evolution

Topological change

Topological change

No Topological change



Cartan's Magic Formula of Continuous Topological Evolution of differential forms

Cartan's Magic Formula

of Continuous Topological Evolution of differential forms

is the

LIE DIFFERENTIAL

acting on a system 1-form of Action, A, with respect to a process direction field, J

Cartan's Magic Formula

$$L_{(J)}A = i(J)dA + d(i(J)A)$$

Cartan's Magic Formula $L_{(J)} A = i(J)dA + d(i(J)A)$

Change notation to yield

$$L_{(J)}A = W + d(U) = Q$$

W = Work 1-form, U = Internal energy, Q = Heat 1-form

Cartan's Magic Formula $L_{(J)} A = i(J)dA + d(i(J)A)$

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W = Work 1-form, U = Internal energy, Q = Heat 1-form

A UNIVERSAL cohomological formulation of the

FIRST LAW of THERMODYNAMICS!!

Kuratowski's Magic Formula

Relative to a Kolmogorov-Cartan T0 topology,

the exterior differential is a

Limit Point generator.

Kuratowski's Magic Formula

Relative to a Kolmogorov-Cartan T0 topology,

the exterior differential is a

Limit Point generator.

For a differential form **\(\Sigma**

Limit Points of $\Sigma = d\Sigma$

This result focuses attention on Cohomology

The Kolmogorov-Cartan Topology is generated by the elements of the Pfaff Sequence of A

Pfaff Sequence :{A, dA, A^dA, dA^dA}

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Pfaff Topological Dimension PTD(A)

PTD(A)= # of non-zero elements in Pfaff Sequence

The Kolmogorov-Cartan Topology is generated by the elements of the Pfaff Sequence of A

Pfaff Sequence :{A, dA, A^dA, dA^dA}

Pfaff Topological Dimension PTD(A)

```
PTD(A) = 1 : \{A, 0, 0, 0\}
```

$$PTD(A) = 2 : \{A, dA, 0, 0\}$$

$$PTD(A) = 3 : \{A, dA, A^dA, 0\}$$

$$PTD(A) = 4 : \{A, dA, A^dA, dA^dA\}$$

The Kolmogorov-Cartan Topology is generated by the elements of the Pfaff Sequence of A

Pfaff Sequence :{A, dA, A^dA, dA^dA}

The Closure of A is the union of A and dA.

The Closure of A^dA is the union of A^dA and dA^dA.

Kolmogorov-Cartan Topology

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Pfaff Sequence :{A, dA, A^dA, dA^dA}

The Closure of A is the union of A and dA.

The Closure of A^dA is the union of A^dA and dA^dA.

The Kolmogorov-Cartan topology has a **Specialization** basis {A, Closure of A, A^dA, Closure of A^dA }.

can describe the irreversible evolution on an

Open non-equilibrium Symplectic domain, PTD 4, with evolutionary orbits being irreversibly attracted to a

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Open non-equilibrium Symplectic domain, PTD 4, with evolutionary orbits being irreversibly attracted to a

Closed non-equilibrium Contact domain, PTD 3, with emergent topological defects (stationary states and coherent structures), and a possible ultimate decay to the

Isolated-Equilibrium Caratheodory (integrable) domain of PTD 2 or less.

 Topological change is a necessary condition for thermodynamic irreversibility.

- 1. Topological change is a necessary condition for thermodynamic irreversibility.
- 2. Continuous non-homeomorphic processes of topological evolution establish a logical basis for thermodynamic irreversibility and the arrow of time without the use of statistics.

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Continuous Processes can represent the evolution from a disconnected topology (≥ 3) to a connected topology (≤ 2).

Continuous Processes can NOT represent the evolution from a connected topology (≤ 2) to a disconnected topology (≥ 3).

Therefore, Connectivity and Continuity determine

A Topological Arrow of Time.

You can describe the decay of turbulence continuously, but NOT the creation of turbulence.

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A Topological Arrow of Time.

You can describe the decay of turbulence continuously, but NOT the creation of turbulence.

Kiehn, R. M. (2003), Thermodynamic Irreversibility and the Arrow of Time, in "The Nature of Time: Geometry, Physics and Perception", R. Bucher et al. (eds.), Kluwer, Dordrecht, Netherlands, 243-250. (http://www22.pair.com/csdc/pdf/arwfinal.pdf)

- 1. Topological change is a necessary condition for thermodynamic irreversibility.
- 2. Continuous non-homeomorphic processes of topological evolution establish a logical basis for thermodynamic irreversibility and the arrow of time without the use of statistics.
- 3. Evolution from a disconnected KCTo topology to a connected topology can be continuous and irreversible, but it is a theorem of topology that a map from a connected topology to a disconnected topology cannot be C2 continuous.

4. C2 Continuous Topological Evolution permits irreversible processes, for which, Q^dQ≠0. Segmented C1 processes approximating smooth C2 processes can be reversible, Q^dQ=0, while the C2 smooth processes are irreversible, Q^dQ≠0.

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- 5. On odd-dimensional spaces, sequential C1 (translational) processes can be thermodynamically reversible, while intransitive C2 processes (rotation and expansion with a fixed point) can be thermodynamically irreversible.

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Engineering Motto for Minimizing energy loss:

Translational Acceleration dot Angular Momentum => 0

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- 5. On odd-dimensional spaces, sequential C1 (translational) processes can be thermodynamically reversible, while intransitive C2 processes (rotation and expansion with a fixed point) can be thermodynamically irreversible.
- 6. The Twin Paradox is resolved if the process paths indicate topological change. Otherwise, there is no disparate aging.

7. Adiabatic processes are transverse to the Heat 1-form, $(i(\rho V_4)Q)=0$. Adiabatic processes need not be quasi-static, and can be reversible or irreversible. Topological Transitions can occur in finite (often short) time.

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Kiehn, R. M. (2008) **Topological Torsion and Macroscopic Spinors**, "Non-Equilibrium Systems and Irreversible Processes Vol 5", Lulu Enterprises, Inc., 3131 RDU Center, Suite 210, Morrisville, NC 27560, see (http://www.lulu.com/kiehn).

- 7. Adiabatic processes are transverse to the Heat 1-form, $(i(\rho V_4)Q)=0$. Adiabatic processes need not be quasi-static, and can be reversible or irreversible. Topological Transitions can occur in finite (often short) time.
- 8. A fundamental difference between Work and Heat is that i(ρV4)W=0, always; but it is not true that i(ρV4)Q=0, always. The Work 1-form, W, is always transverse to the process, ρV4, but the Heat 1-form, Q, may or may not be transverse; the Heat 1-form, Q can have longitudinal components in the direction of the process. Such is the subtle topological difference between Work and Heat.

 For non-equilibrium systems, the 3-form of Topological Torsion (an N-1=3-form current) is not zero:

 $A^dA=i(T_4)dx^dy^dz^dt\neq 0.$

The Topological Torsion vector, T₄, is deduced intrinsically from the 1-form that encodes the thermodynamic system.

It can be used as a direction field for a process current, pT₄.

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The Topological Torsion vector, **T**₄, is deduced intrinsically from the 1-form that encodes the thermodynamic system. It can be used as a direction field for a process current, **ρT**₄.

10. For PTD=3 "closed" thermodynamic systems, the process current has zero divergence, and the 4D volume element is a conformal invariant (any ρ). This result is the space-time extension of the Liouville theorem that preserves the phase-space volume element in classical theory

11. For a PTD=4 "open" thermodynamic systems, the Topological Torsion vector does not have zero divergence, and so the process current ρT4 may not be closed for arbitrary ρ (that is, the divergence of the process current is not zero).

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This result is the extension of the Vlasov equation; the 4D differential volume element is expanding or contracting. Such processes in the direction of T₄ are irreversible and dissipative.

12. A major result is that the Kolmogorov-Cartan T₀ topology is a disconnected topology for non-equilibrium systems (PTD=4,PTD=3) and is a connected topology for equilibrium systems (PTD=2,PTD=1).

- 12. A major result is that the Kolmogorov-Cartan T₀ topology is a disconnected topology for non-equilibrium systems (PTD=4,PTD=3) and is a connected topology for equilibrium systems (PTD=2,PTD=1).
- 13. A key artifact of non-equilibrium is the existence of

Topological Torsion current 3-forms, J_{Torsion},

Topological Spin current 3-forms, J_{Spin},

Topological Adjoint current 3-forms, Jadjoint.

These 3-forms are similar to the Ampere current 3-form, J_{Ampere},

BUT

where $d J_{Ampere} = 0$, always,

the other current 3-forms are not closed unless they are

homogeneous of degree zero.

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where $d J_{Ampere} = 0$, always,

the other current 3-forms are not closed unless they are homogeneous of degree zero.

NOTE: Any 3-form admits (many) integrating factors that will make the 3-form homogenous of degree zero.

The Topological Torsion 3-form is related to Helicity,
The Topological Spin 3-form is related to Spin,

The Adjoint 3-form is related to the interaction energy.

The **Topological Torsion 3-form** is related to **Helicity**,
The **Topological Spin 3-form** is related to **Spin**,
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All three are related to different species of dissipative phenomena, which only occur in non-equilibrium systems.

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The dissipation coefficients are related to the non-zero divergences of the vector coefficients of each 3-form.

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The dissipation coefficients are related to the non-zero divergences of the vector coefficients of each 3-form.

For example, in electromagnetic systems, the dissipation coefficient is proportional to EoB; in hydrodynamics, the dissipation coefficient is called "Bulk viscosity".

14. Examples of thermodynamic systems can be given to demonstrate that the conjectured format of the London Current of superconductivity, where J=χA, can be deduced as an emergent consequence of the Topological Theory of Thermodynamics

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- 15 Examples can generate a Spin Current 3-form, S, where, formally, the Spin current is proportional to the Lorentz force (the space-time components of the Work 1-form, W).

- 14. Examples of thermodynamic systems can be given to demonstrate that the conjectured format of the London Current of superconductivity, where J=χA, can be deduced as an emergent consequence of the Topological Theory of Thermodynamics
- 15 Examples can generate a Spin Current 3-form, where, formally, the Spin current is proportional to the Lorentz force (the space-time components of the Work 1-form, W).

This is a new interpretation of an old result, $J=\sigma(E+VxB)$, which is Ohm's law. The new part is due to the idea that the dissipation is due to Spin Currents and the transport of collective spins, A^{G} .

16. In the PTD=4 case, there exist density distributions, ρ, such that the divergence of the process current is zero.

- 16. In the PTD=4 case, there exist density distributions, ρ , such that the divergence of the process current is zero.
 - There exist an infinite number of such integrating factors, that define "stationary states" far from equilibrium.

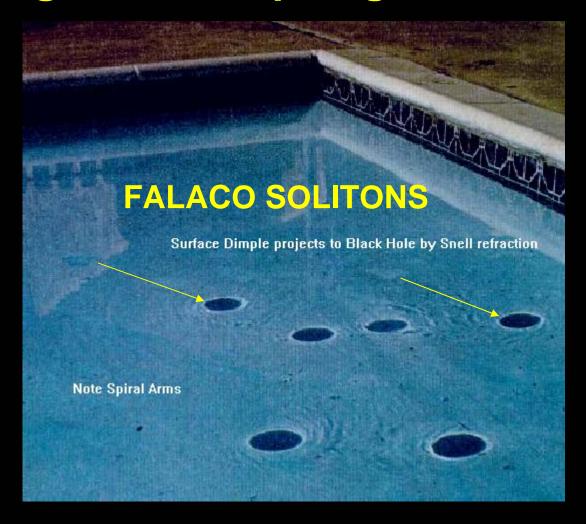
16. In the PTD=4 case, there exist density distributions, ρ, such that the divergence of the process current is zero.

There exist an infinite number of such integrating factors, that define "stationary states" far from equilibrium.

It can be demonstrated in terms of continuous topological evolution that a density distribution which defines a "stationary" state can

emerge as a topological defect.

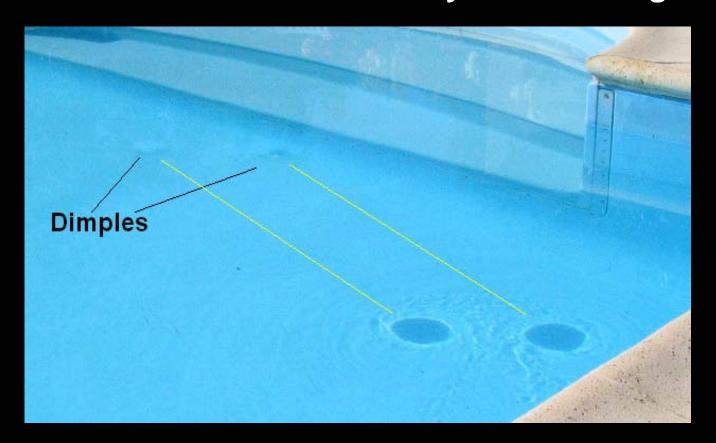
in a PTD=4 system, by means of a dissipative processes.



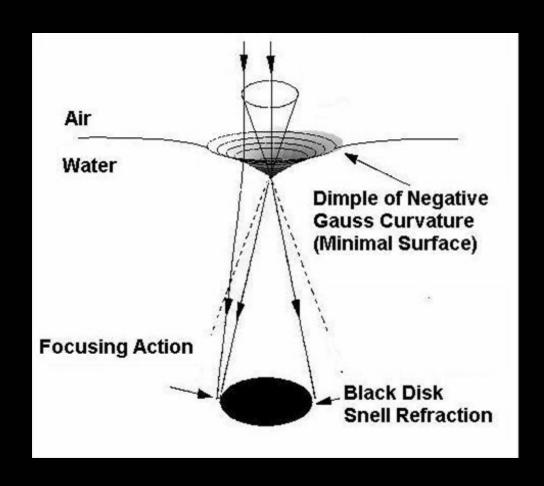
Long Lived Topological Defects in a Swimming Pool

Creation time < 5 seconds. Lifetime > 15 minutes

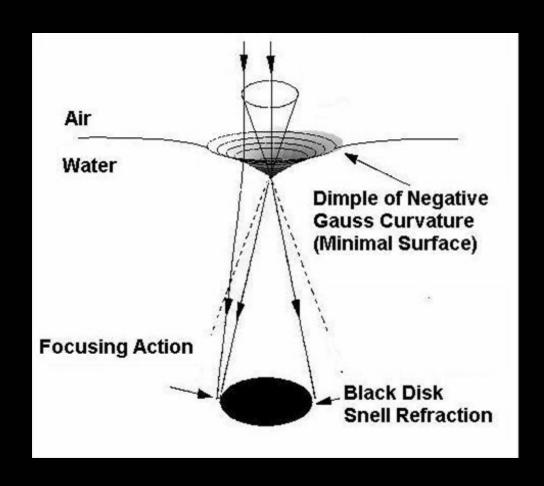
FALACO SOLITONS Movie by D. Radabaugh



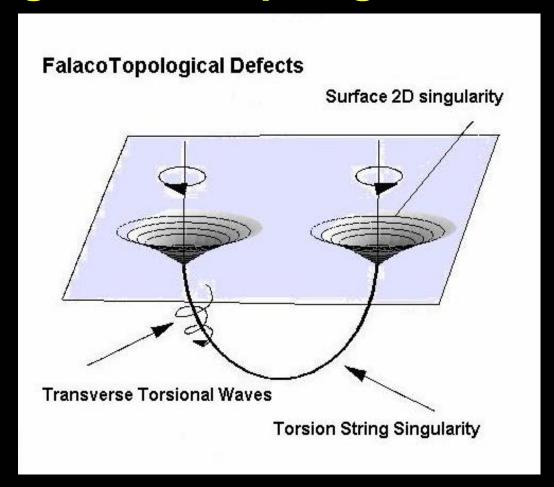
Solar Elevation about 30 degrees (See movie at http://www22.pair.com/csdc/download/spotsmovie.avi



Snell refraction of Falaco Soliton Spin Pairs



Snell refraction of Falaco Soliton Spin Pairs



The first measurableTorsion String coupling between branes

This real world effect has been ignored by string theorists !!!

16. In the PTD=4 case, there exist density distributions, ρ , such that the divergence of the process current is zero.

There exist an infinite number of such integrating factors, that define "stationary states" far from equilibrium.

It can be demonstrated in terms of continuous topological evolution that a density distribution which defines a "stationary" state can emerge as a topological defect in a PTD=4 system, by means of a dissipative processes.

Such a result gives formal credence to Prigogine's conjectures.

17. The topological structure of domains of PTD=3, or more, creates a disconnected topology of multiple components. If solutions to a particular evolutionary problem exist, then the solutions are not unique.

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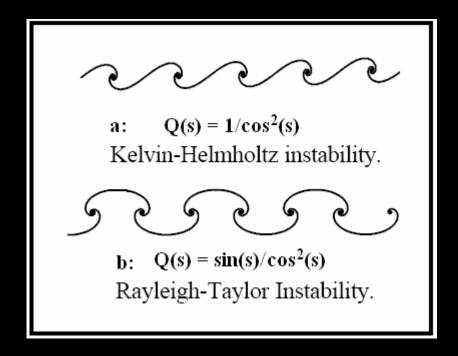
A PTD>2 non-equilibrium thermodynamic system always has a non-zero Topological Torsion 3-form.

17. The topological structure of domains of PTD=3, or more, creates a disconnected topology of multiple components. If solutions to a particular evolutionary problem exist, then the solutions are not unique.

A PTD>2 non-equilibrium thermodynamic system always has a non-zero Topological Torsion 3-form, A^F.

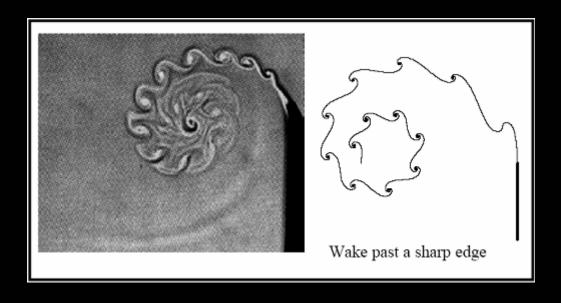
Envelope solutions, such as Huygen wavelets and propagating tangential discontinuities (called signals, or wakes) are classic examples of solution non-uniqueness. Topological Torsion is an artifact of non-uniqueness, and of Turbulence.

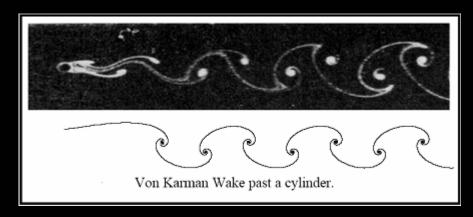
Hydrodynamic Wakes as topological "limit points"



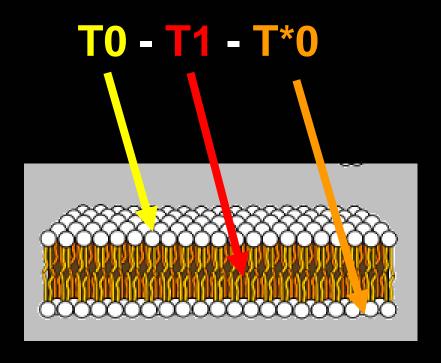
Exact solutions of Hydrodynamic instabilities !!!

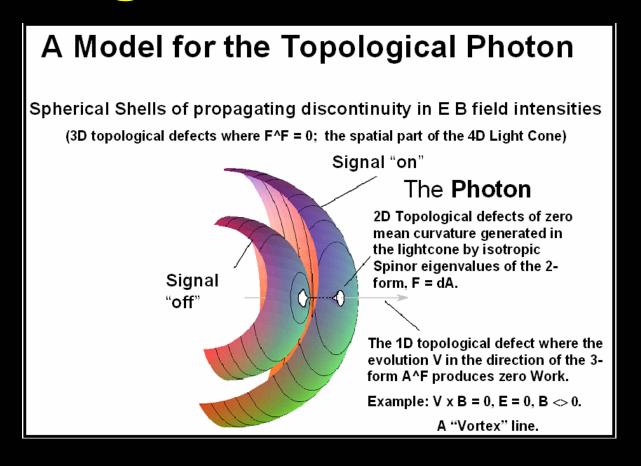
dx/ds = cos(Q), dy/ds = sin(Q)





The Topological double layer Membrane





The EM signal (the Photon) as a Topological double layer Membrane ... T0-T1-T*0

18. All Hamiltonian, Symplectic-Bernoulli and Helmholtz processes are thermodynamically reversible. In particular, the work 1-form, W, created by Hamiltonian processes is of Pfaff Topological Dimension 1 or less. In all reversible cases the Work 1-form is closed, dW=0.

- 18. All Hamiltonian, Symplectic-Bernoulli and Helmholtz processes are thermodynamically reversible. In particular, the work 1-form, **W**, created by Hamiltonian processes is of Pfaff Topological Dimension 1 or less. In all reversible cases the Work 1-form is closed, **dW=0**.
- 19. The assumption of uniqueness of evolutionary solutions (which implies the Pfaff Topological Dimension of the thermodynamic system be equal to 2 or less), or the assumption of homeomorphic evolution, have imposed constraints upon classical mechanics that eliminate any time asymmetry.

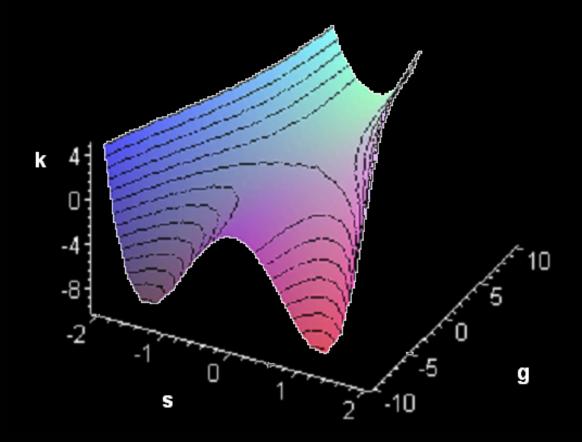
20. The Lie differential acting on differential forms is not necessarily the same as a linear affine covariant differential acting on tensors. If the process is locally adiabatic (no heat flow in the direction of the evolutionary process), then the Lie differential and the covariant differential can be made to coincide, as they both satisfy the Koszul axioms for an affine connection.

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This is a surprising result, for, when the argument is reversed, the theorem implies that the ubiquitous affine covariant differential of tensor analysis, acting on a 1-form of Action, can always be cast into a form representing an adiabatic process. **Warning**: Restrictions of processes which satisfy the constraints of tensor analysis, and use an affine integrable connection to define Covariant derivatives, are always adiabatic.

21. On spaces of PTD=4, the Jacobian of the components of the 1-form of Action, A, define a correlation matrix, which has a characteristic polynomial that defines an equation of state in terms of Cayley-Hamilton similarity invariants.

Universal Topological Thermodynamic Phase Function



A van der Waals gas with a Higgs potential,

An Envelope of a 4D Cayley-Hamilton characteristic polynomial

Universal Topological Thermodynamic Phase Function

The 4D universal topological phase function can be used to explain **Spinodal Decomposition**, and give a topological insight into the Gibbs coexistent phase formula.

Universal Topological Thermodynamic Phase Function

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The 4D universal topological phase function also can be used in dimensions greater than 4 in order to represent multi-component potentials in Chemistry Reactions. The results then can be pulled back to the 4D differential variety of measurement.

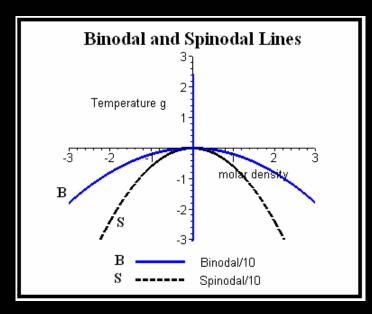
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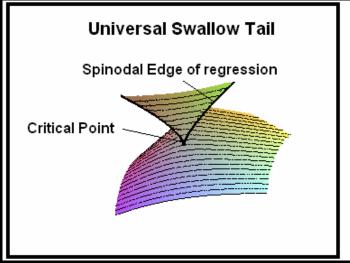
The Cayley-Hamilton theorem produces an implicit hypersurface function that can have envelopes and edges of regression in the format of the Gibbs function for a (universal and deformable) van der Waals gas.

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The method yields analytic expressions for the critical point, and the binodal and spinodal lines, in terms of the similarity invariants. The same technique can be applied to dynamical systems.



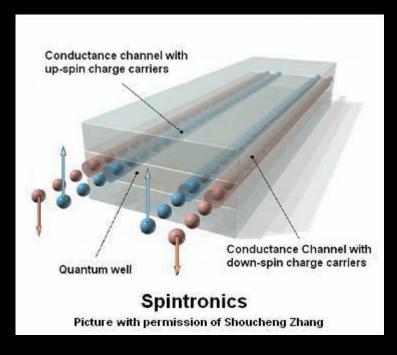


22. Cartan's Magic formula, in terms of the Lie differential acting on exterior differential 1-forms, establishes the long sought for combination of dynamics and thermodynamics, enabling non-equilibrium systems and many irreversible processes to be computed in terms of continuous topological evolution, without resort to probability theory and statistics.

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- 23. **Topological fluctuations** can be induced by processes that have components in terms of thermodynamic macroscopic Spinors. Thermodynamic **Macroscopic Spinors** are non-zero complex eigenvectors with complex eigenvalues (and zero quadratic form) of the antisymmetric 2-form (or matrix) representing the "Limit Points", dA, of the 1-form of Action, A. Such Macroscopic Spinors are capable of representing minimal surface conjugate pairs.

24. **Topological Insulators** correspond to Impedances defined in terms of quantized Topological Spin, and Spin pairs coupled via Falaco Solitons. If the divergence of the Spin 4vector vanishes, the **Spin Current** is time reversal invariant.

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- 25. The fact that any synergetic system of parts in effect defines a topology implies that the universal method of topological thermodynamics is applicable to economic systems, political systems, as well as to biological systems. Such systems admit long-lived states far from equilibrium (adulthood). However, the theory can be used to design subtle perturbations to such systems such that their metastable lifetimes can be extended or destroyed.

26. The thermodynamic processes that lead to self-similarity of a Current 3-form L_(J)C=σ C can generate **fractals** and **holographic** effects where small neighborhoods replicate the whole, approximately. The homogeneity coefficient is the trace of the Jacobian Collineation: σ = Trace[∂C^m/∂xⁿ], or the divergence of the Process vector field.

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Cosmology in terms of a non-equilibrium

Van der Waals Gas explains

Cosmology as a non-equilibrium Van der Waals Gas explains

a.) The granularity of the night sky as exhibited by stars and galaxies due to density fluctuations near the critical point, and the Newtonian law of gravitational attraction proportional to 1/r² as a correlation between fluctuations (due to Lev Landau).

Cosmology as a non-equilibrium Van der Waals Gas explains

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- c.) The possibility of domains of negative pressure (explaining what has recently been called "dark energy") are due to a classical "Higgs" mechanism for aggregates below the critical temperature.

Cosmology as a non-equilibrium Van der Waals Gas explains

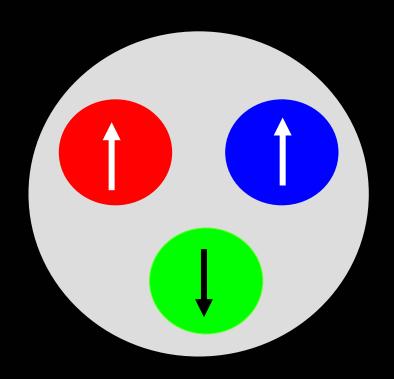
d.) The possibility of domains of negative temperature (explaining what has recently been called "dark matter") are due to macroscopic collective states of ordered spins. The conjecture is that Positive temperature radiates, Negative temperature does not. The conjecture is that black holes could be negative temperature states of collective spins.

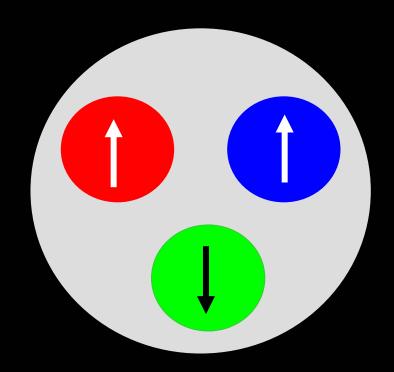
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- e.) The possibility of domains where gravitational effects (quadratic similarity invariants, or 2nd order Gauss curvature effects) appear to be related to entropy and temperature properties of the thermodynamic system, and where cubic curvature effects due to Spin and Adjoint current 3-forms could impede gravitational collapse.

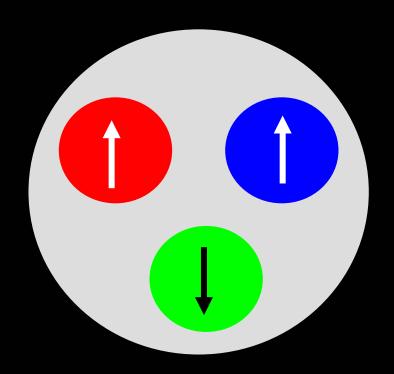
Cosmology as a non-equilibrium Van der Waals Gas explains

f.) Black Holes (generated by Petrov Type D solutions in gravitational theory) are to be related to Minimal Surface solutions to the Universal thermodynamic 4th order Phase function.

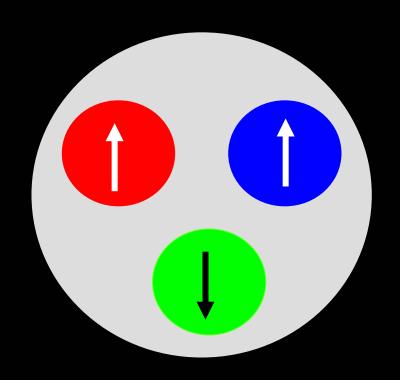




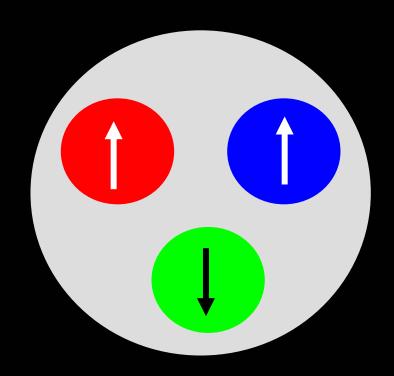
Two Upper sets and One Lower Set?



A specialization preorder system



Or a **PROTON**



Or a **PROTON** Made up from **QUARKS**

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Both Kolomogorov topologies are partitions of a DISCRETE Alexandroff T1 topology. The T1 topology is inherent in the concept of DISCRETE PARTICLES.

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Is this the

Universality of Topological Thermodynamics

acting as the foundation for Quarks?

Thanks for your interest

Contact Professor R. M. Kiehn at

rkiehn2352@ aol.com http://www.cartan.pair.com

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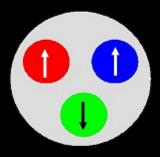
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The Universal Effectiveness of

Topological Thermodynamics

From the Perspective of Continuous Topological Evolution



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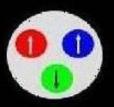


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