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> restart:
Ranada2.mws (RMKiehn 9/10/97)
A Propagating solution with Longitudinal B. By modifying the Ranada stereographic example slightly it is
possible to generate a propagating solution to the Maxwell Faraday equations that is irreducibly 3
dimensional. The Electric field is always transverse to the surface of propagating phase, but the B field is
not.
>
> restart;with(liesymm):with(linalg):with(plots):setup(x,y,z,t);
Warning, new definition for close
Warning, new definition for norm
Warning, new definition for trace
[x, y, z, t]
> defform(x=0,y=0,z=0,t=0,a=const,b=const,c=const,k=const,mu=const,omega=const,m=c
onst);
defform(x=0,y=0,z=0,t=0,a=const,b=const,c=const,k=const,mu=const,omega=const,m=const)
> dR:=[d(x),d(y),d(z),d(t)];
dR := [d(x), d(y), d(z), d(t)]
Define a characteristic denominator, and then specify the four functions that are the covariant components
of the Action 1-form.
> p:=2;e:=0;c:=1;a:=1;b:=1;n:=4;rrxy:=1+a*(x)^p+b*(y)^p+c*(z)^p+e*t^p;ff:=cos(k*z-
omega*t)/rrxy^(n/p);
> ;
p := 2
e := 0
c := 1
a := 1
b := 1
n := 4
ff :=  $\frac{\cos(-kz + \omega t)}{(1 + x^2 + y^2 + z^2)^2}$ 
A choice for the 4 potentials. Note that p has been chosen as 2 to yield the quadratic form, n is
taken to be 4 to fit the Ranada format, the phase propagation is presumed to be in the positive z direction,
and the speed of propagation is omega/k.
> A1:=y*ff;A2:=-x*ff;A3:=k*cos(k*z-omega*t);A4:=omega*cos(k*z-omega*t);
A1 :=  $\frac{y \cos(-kz + \omega t)}{(1 + x^2 + y^2 + z^2)^2}$ 
A2 :=  $-\frac{x \cos(-kz + \omega t)}{(1 + x^2 + y^2 + z^2)^2}$ 
A3 :=  $k \cos(-kz + \omega t)$ 
A4 :=  $\omega \cos(-kz + \omega t)$ 
> Action:=A1*d(x)+A2*d(y)+A3*d(z)-A4*d(t);
Action :=  $\frac{y \%1 d(x)}{(1 + x^2 + y^2 + z^2)^2} - \frac{x \%1 d(y)}{(1 + x^2 + y^2 + z^2)^2} + k \%1 d(z) - \omega \%1 d(t)$ 
%1 :=  $\cos(-kz + \omega t)$ 
> F:=wcollect(d(Action));

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$$F := \frac{y(-k\%1 - k\%1 x^2 - k\%1 y^2 - k\%1 z^2 + 4\%2 z)((d(x)) \&^{\wedge}(d(z)))}{(1+x^2+y^2+z^2)^3} + \left(-\frac{\%2(1-3x^2+y^2+z^2)}{(1+x^2+y^2+z^2)^3} - \frac{\%2(1+x^2-3y^2+z^2)}{(1+x^2+y^2+z^2)^3} \right) ((d(x)) \&^{\wedge}(d(y))) - \frac{x(-k\%1 - k\%1 x^2 - k\%1 y^2 - k\%1 z^2 + 4\%2 z)((d(y)) \&^{\wedge}(d(z)))}{(1+x^2+y^2+z^2)^3} - \frac{x\%1 \omega((d(y)) \&^{\wedge}(d(t)))}{(1+x^2+y^2+z^2)^2} + \frac{y\%1 \omega((d(x)) \&^{\wedge}(d(t)))}{(1+x^2+y^2+z^2)^2}$$

%1 := sin(-k z + ω t)
%2 := cos(-k z + ω t)

F is the electromagnetic 2-form in covariant language for all diffeomorphisms. It is gauge invariant with respect to all closed 1-form additions to the 1-form of Action.

The three components of the Vector potential are:

> **A:=[A1,A2,A3];**

$$A := \left[\frac{y \cos(-k z + \omega t)}{(1+x^2+y^2+z^2)^2}, -\frac{x \cos(-k z + \omega t)}{(1+x^2+y^2+z^2)^2}, k \cos(-k z + \omega t) \right]$$

The three components of the Magnetic field are:

> **B:=(curl(A,[x,y,z]));B1:=(B[1]);B2:=(B[2]);B3:=(B[3]);**

$$B1 := \frac{x \sin(-k z + \omega t) k}{(1+x^2+y^2+z^2)^2} - 4 \frac{x \cos(-k z + \omega t) z}{(1+x^2+y^2+z^2)^3}$$

$$B2 := \frac{y \sin(-k z + \omega t) k}{(1+x^2+y^2+z^2)^2} - 4 \frac{y \cos(-k z + \omega t) z}{(1+x^2+y^2+z^2)^3}$$

$$B3 := -2 \frac{\cos(-k z + \omega t)}{(1+x^2+y^2+z^2)^2} + 4 \frac{x^2 \cos(-k z + \omega t)}{(1+x^2+y^2+z^2)^3} + 4 \frac{y^2 \cos(-k z + \omega t)}{(1+x^2+y^2+z^2)^3}$$

Note that B3 is not zero and is perpendicular to the propagating phase front.

The three components of the Electric Field are:

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> **E:=[-diff(A4,x)-diff(A[1],t),-diff(A4,y)-diff(A[2],t),-diff(A4,z)-diff(A[3],t)];**
E1:=factor(E[1]);E2:=factor(E[2]);E3:=factor(E[3]);

$$E1 := \frac{y \sin(-k z + \omega t) \omega}{(1+x^2+y^2+z^2)^2}$$

$$E2 := -\frac{x \sin(-k z + \omega t) \omega}{(1+x^2+y^2+z^2)^2}$$

$$E3 := 0$$

It would appear that this "wave" is TE (E in the direction of advancing phase is zero.)

The Topological Parity 4 form (Second Poncare invariant) vanishes implying that the 1 form of Action defines a contact manifold. Therefore the integral of the 3-form of topological torsion over a closed domain, not a boundary, is a topological invariant.

> **EdotB:=factor(innerprod(E,B));EdotE:=innerprod(E,E);**

$$EdotB := 0$$

$$EdotE := \frac{\sin(-kz + \omega t)^2 \omega^2 (y^2 + x^2)}{(1 + x^2 + y^2 + z^2)^4}$$

The Torsion current.

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> ExA:=crossprod(E,A):Bphi:=[B1*A4,B2*A4,B3*A4]:
> TORS:=evalm(ExA+A4*B):
> AdotB:=factor(inner(A,B));
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$$AdotB := -2 \frac{k \cos(-kz + \omega t)^2 (1 - x^2 - y^2 + z^2)}{(1 + x^2 + y^2 + z^2)^3}$$

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> TORSION:=[factor(TORS[1]),factor(TORS[2]),factor(TORS[3]),AdotB];
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$$TORSION := \left[-4 \frac{x \omega \%1^2 z}{(1 + x^2 + y^2 + z^2)^3}, -4 \frac{\omega y \%1^2 z}{(1 + x^2 + y^2 + z^2)^3}, -2 \frac{\%1^2 (1 - x^2 - y^2 + z^2) \omega}{(1 + x^2 + y^2 + z^2)^3}, -2 \frac{k \%1^2 (1 - x^2 - y^2 + z^2)}{(1 + x^2 + y^2 + z^2)^3} \right]$$

%1 := cos(-kz + \omega t)

The 4 vector of Torsion Current (and the fourth component which is the helicity) are proportional to the SQUARE of the oscillatory term phase term.

The 4Divergence of the Torsion current is:

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> DIV4DofT:=factor(diverge(TORSION,[x,y,z,t]));
DIV4DofT := 0
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establishing the "conservation law" for topological torsion.

Note that none of this assumed a group constraint, nor an assumption of "duality" to produce an E field from a B field (as used by Ranada).

These fields satisfy the Maxwell-Faraday equations for E and B, by construction!

No mention is made of D and H, nor of the charge distributions and constitutive relations that might be compatible with these field intensities. However, as the Torsion Current is closed, it is a candidate for the charge current distribution of classical electromagnetic theory. The problem is to find solutions to the partial differential equations, curlH-partialdD/dt = T and divD = A.B, where the RHS is given. Another day.

RMK.