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> restart:
The rolling ball with friction.
ROLLBALL.MWS
> with(liessymm):with(linalg):with(plots):with(diffforms):
>
Warning, new definition for close
Warning, new definition for norm
Warning, new definition for trace
Warning, new definition for `&^`
Warning, new definition for d
Warning, new definition for mixpar
Warning, new definition for wdegree

> setup(x,theta,t,v,omega,s);
The independent variables are on a 6 dimensional space
[x, theta, t, v, omega, s]
>
> deform(theta=0,omega=0,x=0,v=0,s=0,t=0,a=const,b=const,c=const,k=const,mu=const
,m=const,beta=const,lambda=const):
> Lagrange:=(1/2*beta*m*lambda^2*omega^2-1/2*m*v^2);
Lagrange :=  $\frac{1}{2} \beta m \lambda^2 \omega^2 - \frac{1}{2} m v^2$ 
Note that the Lagrangian is defined as the Difference between the rotational energy and the translational
energy. !!!!
> Ptheta:=evalm(diff(Lagrange,omega));Px:=evalm(diff(Lagrange,v));Ham:=Ptheta*omeg
a+Px*v-Lagrange;
Ptheta :=  $\beta m \lambda^2 \omega$ 
Px :=  $-m v$ 
Ham :=  $\frac{1}{2} \beta m \lambda^2 \omega^2 - \frac{1}{2} m v^2$ 
Momenta defined canonically,
> Action:=wcollect(Ptheta*d(theta)+Px*d(x)-Ham*d(t)+s*(lambda*d(theta)-d(x)));
The Action is defined in terms of the rotational and potential energy. The momenta are presumed to be
canonical Ptheta=beta*m*lambda^2*omega, Px=m*v. The constraint of no slip is lambda*d(theta)-d(x)
and its Lagrange multiplier is s. The Lagrange function is defined as
L=1/2*beta*m*lambda^2*omega^2+1/2*m*v^2)
Action :=  $(-m v - s) d(x) + (\beta m \lambda^2 \omega + s \lambda) d(\theta) + \left(-\frac{1}{2} \beta m \lambda^2 \omega^2 + \frac{1}{2} m v^2\right) d(t)$ 
> DA:=d(Action);
DA :=  $-m ((d(v)) \wedge (d(x))) - ((d(s)) \wedge (d(x))) + \beta m \lambda^2 ((d(\omega)) \wedge (d(\theta))) + \lambda ((d(s)) \wedge (d(\theta)))$ 
 $- \beta m \lambda^2 \omega ((d(\omega)) \wedge (d(t))) + m v ((d(v)) \wedge (d(t)))$ 
The two form is of rank 3 hence when the no slip condition fails. Hence the two form defines a non-compact
symplectic space.
> Hel:=Action&^DA;
Hel :=  $\left(-\frac{1}{2} m^2 v^2 - m v s - \frac{1}{2} m^2 \beta \lambda^2 \omega^2\right) \wedge (d(x), d(v), d(t)) + (m v + s) \beta m \lambda^2 \omega \wedge (d(x), d(\omega), d(t))$ 
 $- (m v + s) \beta m \lambda^2 \wedge (d(x), d(\omega), d(\theta)) + (-\lambda m v + \beta m \lambda^2 \omega) \wedge (d(x), d(s), d(\theta))$ 

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$$\begin{aligned}
& + \lambda (\beta m \lambda \omega + s) m v \wedge (d(\theta), d(v), d(t)) + \left( -\frac{1}{2} \beta^2 m^2 \lambda^4 \omega^2 - \beta m \lambda^3 \omega s - \frac{1}{2} m^2 \beta \lambda^2 v^2 \right) \wedge (d(\theta), d(\omega), d(t)) \\
& - \lambda (\beta m \lambda \omega + s) m \wedge (d(\theta), d(v), d(x)) - \frac{1}{2} m (\beta \lambda^2 \omega^2 - v^2) \lambda \wedge (d(t), d(s), d(\theta)) \\
& + \left( \frac{1}{2} \beta m \lambda^2 \omega^2 - \frac{1}{2} m v^2 \right) \wedge (d(t), d(s), d(x))
\end{aligned}$$

The Helicity 3-form above does not vanish

> **DADA:=wcollect((DA^DA));**

$$\begin{aligned}
DADA := & 2 m^2 \beta \lambda^2 \omega \wedge (d(v), d(x), d(\omega), d(t)) - 2 m^2 \beta \lambda^2 \wedge (d(v), d(x), d(\omega), d(\theta)) \\
& - 2 m \lambda \wedge (d(v), d(x), d(s), d(\theta)) + 2 \beta m \lambda^2 \omega \wedge (d(s), d(x), d(\omega), d(t)) - 2 m v \wedge (d(s), d(x), d(v), d(t)) \\
& - 2 \lambda^3 \beta m \omega \wedge (d(s), d(\theta), d(\omega), d(t)) + 2 \beta m^2 \lambda^2 v \wedge (d(\omega), d(\theta), d(v), d(t)) \\
& + 2 \lambda m v \wedge (d(s), d(\theta), d(v), d(t)) - 2 \beta m \lambda^2 \wedge (d(s), d(x), d(\omega), d(\theta))
\end{aligned}$$

> **Torsion:=wcollect(factor(Hel^DA));**

$$\begin{aligned}
Torsion := & m^2 \lambda (-v^2 - \beta \lambda^2 \omega^2 + 2 \lambda v \beta \omega) \wedge (d(x), d(v), d(t), d(s), d(\theta)) \\
& + m^2 \lambda (m \beta^2 \omega^2 \lambda^3 - m v^2 \beta \lambda - 2 \lambda \beta v s + 2 \lambda^2 \beta \omega s) \wedge (d(x), d(v), d(t), d(\omega), d(\theta)) \\
& + m^2 \lambda (-\beta^2 \omega^2 \lambda^3 + 2 v \beta \lambda^2 \omega - v^2 \beta \lambda) \wedge (d(x), d(\omega), d(t), d(s), d(\theta))
\end{aligned}$$

> **Vol:=factor(DA^DA^DA);**

$$Vol := 6 m^2 \beta \lambda^2 (-v + \lambda \omega) \wedge (d(v), d(x), d(\omega), d(t), d(s), d(\theta))$$

> **scale:=scalarpart(factor(DA^DA^DA));**

$$scale := 6 m^2 \beta \lambda^2 (-v + \lambda \omega)$$

>

The Torsion 5-form has 6 components but 3 of them are zero.

> **Tv:=- (scalarpart(Torsion^d(v))\*3/scale);**

$$Tv := -\frac{1}{2} \frac{-\beta^2 \omega^2 \lambda^3 + 2 v \beta \lambda^2 \omega - v^2 \beta \lambda}{\lambda \beta (-v + \lambda \omega)}$$

> **To:=- (scalarpart(Torsion^d(omega))\*3/scale);**

$$To := -\frac{1}{2} \frac{-v^2 - \beta \lambda^2 \omega^2 + 2 \lambda v \beta \omega}{\lambda \beta (-v + \lambda \omega)}$$

> **Ts:=factor(scalarpart(Torsion^d(s))\*3/scale);**

$$Ts := \frac{1}{2} \frac{\beta m \lambda^2 \omega^2 - m v^2 - 2 v s + 2 \lambda \omega s}{-v + \lambda \omega}$$

> **SS:=solve(Ts,s);**

$$SS := -\frac{1}{2} \frac{m (\beta \lambda^2 \omega^2 - v^2)}{-v + \lambda \omega}$$

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The Torsion 5-form has 3 non-zero components on the six dimensional space and 3 zero components.. The non-zero components are orthogonal to the spatial coordinates {x,theta,t} in the 6D space{x,theta,t,v,omega,s}

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The above 6-form is the volume element of the 6 dimensional space. Note that the volume vanishes at the no slip condition. This 6-dimensional space is the top Pfaffian. It is NOT compact without boundary.

> **AT:=factor(Hel^DADA);**

As this 7 form vanishes, it is proved that the Pfaff dimension of the Action 1-form is 6, with the coordinates  $\{x, \theta, t, v, \omega, s\}$

Note that the 6-form vanishes when the noslip condition prevails. The space then reduces to a contact manifold of 5 dimensions. However, initial conditions do not necessarily satisfy the no slip constraint. Therefore the system decays (irreversibly) until the no slip condition is satisfied. Thereafter the dynamics is that of a Hamiltonian system with no irreversible dissipation.

$$AT := 0$$

>

> **AVEC := [Px-s, Ptheta+s\*lambda, -Ham, 0, 0, 0];**

$$AVEC := \left[ -m v - s, \beta m \lambda^2 \omega + s \lambda, -\frac{1}{2} \beta m \lambda^2 \omega^2 + \frac{1}{2} m v^2, 0, 0, 0 \right]$$

> **TORS := [0, 0, 0, Tv, To, Ts];**

**TORS :=**

$$\left[ 0, 0, 0, -\frac{1 - \beta^2 \omega^2 \lambda^3 + 2 v \beta \lambda^2 \omega - v^2 \beta \lambda}{\lambda \beta (-v + \lambda \omega)}, -\frac{1 - v^2 - \beta \lambda^2 \omega^2 + 2 \lambda v \beta \omega}{\lambda \beta (-v + \lambda \omega)}, \frac{1}{2} \frac{\beta m \lambda^2 \omega^2 - m v^2 - 2 v s + 2 \lambda \omega s}{-v + \lambda \omega} \right]$$

> **test := innerprod(TORS, AVEC);**

$$test := 0$$

> **DA;**

$$-m ((d(v)) \wedge (d(x))) - ((d(s)) \wedge (d(x))) + \beta m \lambda^2 ((d(\omega)) \wedge (d(\theta))) + \lambda ((d(s)) \wedge (d(\theta))) - \beta m \lambda^2 \omega ((d(\omega)) \wedge (d(t))) + m v ((d(v)) \wedge (d(t)))$$

> **WORK := wcollect(factor((m\*Tv\*(d(x)-v\*d(t))+beta\*m\*lambda^2\*To\*(d(theta))-omega\*d(t))+Ts\*(lambda\*d(theta)-d(x)))));**

$$WORK := (-m v - s) d(x) + (\beta m \lambda^2 \omega + s \lambda) d(\theta) + \left( -\frac{1}{2} \beta m \lambda^2 \omega^2 + \frac{1}{2} m v^2 \right) d(t)$$

> **factor(Action&^WORK);**

> **Action;**

$$0$$

$$(-m v - s) d(x) + (\beta m \lambda^2 \omega + s \lambda) d(\theta) + \left( -\frac{1}{2} \beta m \lambda^2 \omega^2 + \frac{1}{2} m v^2 \right) d(t)$$

The virtual work is proportional to the Action. The virtual work 1-form is not zero until the no-slip condition is obtained.

> **Wx := factor(scalarpart(WORK&^d(theta)&^d(t)));**

>

$$Wx := -m v - s$$

> **Wo := factor(simplify(scalarpart(WORK&^d(t)&^d(x))));**

$$Wo := \lambda (\beta m \lambda \omega + s)$$

> **Wt := factor(scalarpart(WORK&^d(x)&^d(theta)));**

$$Wt := -\frac{1}{2} m (\beta \lambda^2 \omega^2 - v^2)$$

The Work 1-form relative to the Torsion vector is proportional to the Action with the factor  $2 * \lambda^3 * m^2 * \beta$

> **DWORK := wcollect(simpform(d(WORK)));**

$$DWORK := -m ((d(v)) \wedge (d(x))) - ((d(s)) \wedge (d(x))) + \beta m \lambda^2 ((d(\omega)) \wedge (d(\theta))) + \lambda ((d(s)) \wedge (d(\theta))) - \beta m \lambda^2 \omega ((d(\omega)) \wedge (d(t))) + m v ((d(v)) \wedge (d(t)))$$

> **IRR:=wcollect(WORK&^DWORK);**

$$\begin{aligned}
 IRR := & \left( -\frac{1}{2} m^2 v^2 - m v s - \frac{1}{2} m^2 \beta \lambda^2 \omega^2 \right) \wedge (d(x), d(v), d(t)) + (m v + s) \beta m \lambda^2 \omega \wedge (d(x), d(\omega), d(t)) \\
 & - (m v + s) \beta m \lambda^2 \wedge (d(x), d(\omega), d(\theta)) + (-\lambda m v + \beta m \lambda^2 \omega) \wedge (d(x), d(s), d(\theta)) \\
 & + \lambda (\beta m \lambda \omega + s) m v \wedge (d(\theta), d(v), d(t)) + \left( -\frac{1}{2} \beta^2 m^2 \lambda^4 \omega^2 - \beta m \lambda^3 \omega s - \frac{1}{2} m^2 \beta \lambda^2 v^2 \right) \wedge (d(\theta), d(\omega), d(t)) \\
 & - \lambda (\beta m \lambda \omega + s) m \wedge (d(\theta), d(v), d(x)) - \frac{1}{2} m (\beta \lambda^2 \omega^2 - v^2) \lambda \wedge (d(t), d(s), d(\theta)) \\
 & + \left( \frac{1}{2} \beta m \lambda^2 \omega^2 - \frac{1}{2} m v^2 \right) \wedge (d(t), d(s), d(x))
 \end{aligned}$$

> **W:=(factor(WORK));**

$$W := -d(x) m v - d(x) s + d(\theta) \beta m \lambda^2 \omega + d(\theta) s \lambda - \frac{1}{2} d(t) \beta m \lambda^2 \omega^2 + \frac{1}{2} d(t) m v^2$$

The Work 1-form relative to the unique Torsion vector is not integrable until the no-slip condition is reached. Hence the Heat 1-form Q is not integrable. Therefor a process defined by the unique Torsion vector is irreversible, until the no slip condition is obtained.

The dynamical system defined by the Torsion vector is

> **d(v)-Tv\*d(tau)=0;d(omega)-To\*d(tau)=0;d(s)-Ts\*d(tau)=0;d(x)-v\*d(t);d(theta)-omeg  
a\*d(t)=0;d(tau)-aa\*d(t)=0;**

$$\begin{aligned}
 d(v) + \frac{1}{2} \frac{(-\beta^2 \omega^2 \lambda^3 + 2 v \beta \lambda^2 \omega - v^2 \beta \lambda) d(\tau)}{\lambda \beta (-v + \lambda \omega)} &= 0 \\
 d(\omega) + \frac{1}{2} \frac{(-v^2 - \beta \lambda^2 \omega^2 + 2 \lambda v \beta \omega) d(\tau)}{\lambda \beta (-v + \lambda \omega)} &= 0 \\
 d(s) - \frac{1}{2} \frac{(\beta m \lambda^2 \omega^2 - m v^2 - 2 v s + 2 \lambda \omega s) d(\tau)}{-v + \lambda \omega} &= 0 \\
 d(x) - v d(t) &= 0 \\
 d(\theta) - \omega d(t) &= 0 \\
 d(\tau) - aa d(t) &= 0
 \end{aligned}$$

> **PH:=To\*d(v)-Tv\*d(omega)=0;**

$$PH := -\frac{1}{2} \frac{(-v^2 - \beta \lambda^2 \omega^2 + 2 \lambda v \beta \omega) d(v)}{\lambda \beta (-v + \lambda \omega)} + \frac{1}{2} \frac{(-\beta^2 \omega^2 \lambda^3 + 2 v \beta \lambda^2 \omega - v^2 \beta \lambda) d(\omega)}{\lambda \beta (-v + \lambda \omega)} = 0$$

> **dPxdt = factor(m\*Tv-Ts);Action;**

$$\begin{aligned}
 dPxdt &= -m v - s \\
 (-m v - s) d(x) + (\beta m \lambda^2 \omega + s \lambda) d(\theta) + \left( -\frac{1}{2} \beta m \lambda^2 \omega^2 + \frac{1}{2} m v^2 \right) d(t) &= 0
 \end{aligned}$$

> **dPthetadt:=factor(beta\*m\*lambda^2\*To+Ts\*lambda);**

>

$$dPthetadt := \lambda (\beta m \lambda \omega + s)$$

>

Both the (New) translational momenta and angular momenta decay at the same exponential rate!!!

[ New Momenta  $(mv+s)$  ,  $(2\beta*\lambda^2*m*\omega+s*\omega)$  are NOT conserved.

[ >

[ The first two equations admit an immediate integral during the dissipative motion along the torsion vector.

[ >

$$v + \lambda \beta \omega = (1 + \beta) V_{noslip}$$

[ >

[ >

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[ The no-slip condition states that  $v=\lambda*\omega$ , so the constant of integration is  $1+\beta$ )